1 Magnitude Scales for Marsquakes Calibrated from InSight Data

- Böse, M.\*, S. Stähler, N. Deichmann, D. Giardini, J. Clinton, P. Lognonné, S. Ceylan, M. van Driel, C.
  Charalambous, N. Dahmen, A. Horleston, T. Kawamura, A. Khan, M. Knapmeyer, G. OrhandMainsant, J.-R. Scholz, F. Euchner, B.W. Banerdt
- 5
- 6 \*Corresponding author: Maren Böse, <u>maren.boese@sed.ethz.ch</u>, Institute of Geophysics, ETH

7 Zurich, Sonneggstrasse 5, 8092 Zurich, Switzerland

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#### 10 ABSTRACT

11 In preparation for the NASA InSight (Interior exploration using Seismic Investigations, Geodesy 12 and Heat Transport) Discovery Program mission, Böse et al. (2018) calibrated magnitude scales 13 for marsquakes that incorporated the pre-launch knowledge on Mars' interior structure and the 14 expected ambient and instrumental noise. Now, using data collected during the first two years 15 after the successful deployment of the InSight very-broadband seismometer on the Martian 16 surface, we revise these relations to account for the seismic and noise characteristics observed on 17 Mars. The data collected so far (until October 12, 2020) includes 485 seismic event detections and 18 suggest that (1) marsquakes are characterized by energy between  $\sim 0.1$  and 10 Hz; (2) though first 19 arriving P- and S-wave phases are regularly identified and assigned, both surface-waves and 20 secondary phase arrivals are extremely challenging to identify; (3) the majority of identified events 21 include a strong excitation of an unexpected 2.4 Hz ground resonance; and (4) so-called high-22 frequency events exist that are visible mainly as guided  $P_g/S_g$  wave trains. In view of these 23 observations, we update our scaling relations for the spectral and body-wave magnitudes,  $M_{W,spec}^{Ma}$ ,  $m_{\rm b}^{\rm Ma}$  and  $m_{\rm bS}^{\rm Ma}$ , and introduce a new magnitude scale,  $M_{2.4}^{\rm Ma}$ , for high-frequency events. We use these 24 25 scales to determine that the magnitudes of events in the current InSight v.5 catalog range between 26 1.1 and 3.7, with event-specific uncertainties  $\sigma_M$  ranging from 0.2 to 0.5. Due to the currently 27 unclear interpretation of high-frequency events, magnitude estimates for these events primarily 28 serve for a relative comparison.

### **30 INTRODUCTION**

31 The NASA InSight (Interior exploration using Seismic Investigations, Geodesy and Heat Transport) 32 Discovery Program mission successfully launched on May 5, 2018, from Vandenberg Air Force 33 Base, California, and landed on Mars about six months later (Banerdt et al., 2020). The deployment 34 of the three-component ultra-sensitive very-broadband (VBB) seismometer (Lognonné et al., 35 2019; 2020) was completed by the end of February 2019 (Giardini et al., 2020). Since then, the 36 Marsquake Service (MQS), as part of the InSight Science Team, routinely inspects manually the 37 continuous 20 samples-per-second (sps) seismic data returned by InSight (Clinton et al., 2018). 38 Event data at higher sampling (100 sps) is recovered whenever possible. (The available bandwidth 39 depends on multiple factors related to satellites around Mars that relay the data, as well as 40 competing data volumes from other Martian missions and even other instruments operating on 41 InSight.) Until October 12, 2020 (first Martian year of the InSight mission), MQS has identified 485 42 seismic events that have been interpreted as marsquakes (InSight Marsquake Service, 2021). All 43 of these events are characterized by energy between  $\sim 0.1 - 10$  Hz and have durations typically 44 ranging from 5 - 30 minutes. A small number of events include energy of up to 35 Hz. Waveform 45 amplitudes typically rise only slightly above the lowest background noise, and this lowest noise 46 typically occurs only for a few hours per sol (a Martian sol is 24h 39.5' long). A detailed description 47 of the event catalog and general characteristics of the data collected so far are given in Clinton et 48 al. (2020), Giardini et al. (2020), and Ceylan et al. (2020); noise and event amplitudes are discussed 49 in Charalambous et al. (2021).

50

The seismic events identified in the InSight MQS catalog fall into two families (Clinton *et al.*, 2020). A low-frequency family, which is similar to terrestrial earthquakes, comprises events with clearly defined P- and S-wave arrivals, even though they may have low signal-to-noise ratios (SNR). A second family of high-frequency events is characterized by long-duration wave trains, similar to guided Po/So waves in the oceanic crust (Kennett & Furumura, 2013). The MQS further divides the two events families into the following five classes:

57 • Low-frequency family:

58 o *low-frequency (LF)* events have energy on all 3 components below 2.4 Hz;

59 o *broadband (BB)* events have energy on 3 components, predominantly below 2.4 Hz, but

also including excitement at and possibly also above 2.4 Hz;

• High-frequency family:

- *high-frequency (HF)* events have energy on 3 components, predominantly at 2.4 Hz and
  above;
- o very high-frequency (VF) events are a special case of HF events with energy on the
   horizontal components at higher frequencies being significantly larger than on the
   vertical;
- 67 o 2.4Hz resonance (2.4Hz) events have energy on 3 components, centered around 2.4 Hz
  68 (ambient) resonance with very limited excitation above or below (likely small amplitude
  69 *HF* events).
- 70

An additional class of signals, *super-high frequency* (*SF*) events, are also included in the MQS catalog. These events have an order of magnitude shorter durations than the other events and only include energy above 5 Hz. They are considered to be associated with a very local source (Dahmen *et al.*, 2021), but further research is needed to provide a stable interpretation of these signals. *SF* events are not considered further in this study.

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77 The InSight catalog v.5 (until 2020/10/12; sol 668) holds a total of 279 seismic events with 78 distance estimates of any kind (and an additional 206 events without distance estimates or of 79 unclear origin), including 20 *LF* ( $20^{\circ} \le \Delta \le 100^{\circ}$ ), 9 *BB* ( $25^{\circ} < \Delta < 90^{\circ}$ ), 50 *HF* ( $20^{\circ} < \Delta < 30^{\circ}$ ), 21 *VF* 80  $(3^{\circ} < \Delta < 30^{\circ})$ , and 179 2.4*Hz* ( $6^{\circ} < \Delta < 40^{\circ}$ ) events (Table 1). The 2.4*Hz* events are likely weak 81 examples of *HF* events that would not be observable, but for the additional excitation of the 2.4 Hz 82 subsurface resonance rising above the background noise (Clinton et al., 2020). VF events are HF 83 events with clear differences in energy between horizontal and vertical components and a rising 84 displacement spectrum on the horizontal components. Following Giardini et al. (2020) and van 85 Driel *et al.* (2021), our current understanding is that the seismic waveforms of events in the high-86 frequency (HF, VF, 2.4Hz) family are dominated by guided crustal waves in a shallow layer. This 87 could be a regional sediment layer, or in the most extreme case, the whole crust and lithosphere, 88 as observed in the oceanic crust on Earth (Kennett & Furumura, 2013). The assumed propagation 89 path of seismic waves in the two event families is illustrated in Figure 1.

91 The MQS assigns event quality according to the following conventions: (1) High-quality events ('A') 92 show multiple clear and identifiable phases and clear polarisation (i.e. they can be located using 93 the probabilistic approach of Böse *et al.*, 2016); (2) medium quality events ('B') show multiple clear 94 and identifiable phases, but no polarisation (i.e. only their distance from the lander can be 95 determined, but no absolute location); (3) low-quality events ('C') show no clear phases, but 96 signals are clearly observed; and (4) suspicious events ('D') show only weakly observed signals or 97 signals that may not be attributable to a seismic event. In addition to this quality scheme, all events 98 are characterized by an SNR measurement (Giardini et al., 2020; Charalambous et al., 2021).

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100 In the v.5 MQS catalog, distance estimates are available for all quality 'A' and 'B' events, and for the 101 majority of quality 'C' and some quality 'D' events. A complete description of the approach is given 102 in Clinton et al. (2020) and is summarized below. For low-frequency family events the 103 methodology follows Böse et al. (2016). Using large sets of travel time tables derived from a priori 104 velocity models, distances are computed using picks identified in either the time or spectral 105 domain that are interpreted as P- and S- phases. This allowed to constrain 15 of the 43 low-106 frequency family events in distance. Further work using alignment of waveform envelopes allows 107 these distances to be refined and additional events that do not have clear onsets to be assigned 108 distances (see Giardini et al., 2020). In this manner, 27 of the 43 (63%) low-frequency family 109 events in the MOS catalog are assigned distances. In cases where there are differences between 110 distances based on phase picks and alignments, the alignment-based distances are preferred. For 111 events in the high-frequency family, picks are typically identified using a STA-LTA filter targeting 112 the 2.4 Hz resonance. Two phase onsets are normally visible, that are interpreted as Pg and Sg. Since 113 the propagation velocities of those are not known, we cannot confidently determine absolute 114 distances. In order to provide an approximate distance, as well as a robust estimate of the relative 115 distances within the high-frequency family, we assume that  $P_g$  propagates with  $v_p$  = 4 km/s and  $S_g$ with  $v_s = v_p/\sqrt{3} = 2.3$  km/s. For 249 of 441 (56%) high-frequency family events in the MQS catalog, 116 117  $P_{\rm g}$  and  $S_{\rm g}$  phases are both identified and we could assign distances this way.

### 119 **Pre-Launch Magnitude Calibrations**

120 Assigning event magnitudes is an important part of characterizing seismicity. In our pre-launch 121 paper (Böse *et al.*, 2018) we had calibrated magnitude scales for marsquakes by simulating the 122 seismic wave propagation through a set of hypothetical 1D Mars models, that incorporated the 123 pre-mission knowledge on Mars interior structure (Smrekar et al., 2019) and the assumed ambient 124 and instrumental noise (Murdoch et al., 2015; Mimoun et al., 2017), using the spectral element 125 solver AxiSEM (Nissen-Meyer et al., 2014; van Driel et al., 2017). Inspired by relations developed for Earth, we had defined six magnitude scales: local Mars magnitude,  $M_{\rm L}^{\rm Ma}$ ; P- and S-wave 126 magnitudes,  $m_{\rm b}^{\rm Ma}$  and  $m_{\rm bS}^{\rm Ma}$ , at a period of 3 s; surface-wave magnitude,  $M_{\rm s}^{\rm Ma}$ ; and spectral 127 magnitudes,  $M_{\rm FB}^{\rm Ma}$  and  $M_{\rm F}^{\rm Ma}$ . We had calibrated these scales relative to the seismic moment 128 129 magnitude at  $M_W = 5.5$  by correcting filtered phase or spectral amplitudes in the simulated 130 waveforms for distance-dependent attenuation using

131

$$M_i^{\text{Ma}} = \log_{10}(A_i) + a_i \log_{10}(\Delta) + c_i \tag{1}$$

133

134 The superscript "Ma" stands for "Mars" (following BSSA accepted standard for planetary 135 magnitude in accordance with IAU practice);  $a_i$  and  $c_i$  are scaling coefficients for the different 136 magnitude types;  $A_i$  is a measure of amplitude that is also dependent on magnitude type – and can 137 be the filtered peak displacement amplitude [m] determined from the waveform time-series 138 within a certain time-window, or the spectral amplitude  $[m/\sqrt{Hz}]$  determined from the square-139 root of the long-period plateau of the power spectral density (PSD) of the displacement;  $\Delta$  is the 140 epicentral distance in degrees [<sup>o</sup>]. For both the local and the surface-wave based magnitude scales, 141 Böse et al. (2018) included an additional term for source depth. In a second regression, Böse et al. 142 (2018) had calibrated the six scales with the known seismic moment,  $M_0$ , of the simulations across 143 the entire magnitude range  $1.0 \le M_W \le 7.0$  to obtain a conversion relation to moment magnitude 144  $M_W$ .

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146 The InSight data collected so far suggest that (1) two families of marsquakes exist (low-frequency 147 and high-frequency; see above), which need separate magnitude scales, (2) the frequency content 148 of the more classic low-frequency events ranges from 0.1 to 1 Hz, the high-frequency events from 149 1.5 to up to 10 (and possibly larger) Hz, (3) though first arriving P- and S-wave phases are regularly 150 identified and assigned, both surface-waves and secondary phase arrivals are extremely 151 challenging to identify and cannot be used for regular analysis, and (4) many high-frequency 152 events are primarily visible as an enhanced excitation of the natural 2.4 Hz ambient mode, which 153 is believed to originate from the substructure below the lander (Ceylan et al., 2020; Dahmen, 154 Zenhäuser et al., subm.). In view of these observations, we update our scaling relations (Böse et al., 155 2018) for the body-wave magnitudes,  $m_{\rm b}^{\rm Ma}$  and  $m_{\rm bS}^{\rm Ma}$ , and introduce a new magnitude scale,  $M_{2.4}^{\rm Ma}$ , 156 specifically, for high-frequency events. The surface-wave magnitude,  $M_{\rm S}^{\rm Ma}$ , is not updated, given 157 the lack of clear surface-wave observations (which is likely due to the small sizes of marsquakes 158 detected so far; Clinton et al., 2020), and as explained in the following section, we replace the 159 spectral magnitude  $M_{FB}^{Ma}$  introduced in Böse *et al.* (2018) by  $M_{W,spec}^{Ma}$ .

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161

### 162 **METHOD**

163 Based on the data in the InSight v.5 event catalog (until 2020/10/12; InSight Marsquake Service, 164 2021) and formulations of Böse et al. (2018), we calibrate in this study four magnitude scales: 165 spectral magnitude, M<sub>W,spec</sub> (computed over the full body-wave train), time-domain P- and Sbody-wave magnitudes,  $m_{\rm b}^{\rm Ma}$  and  $m_{\rm bS}^{\rm Ma}$ , and 2.4 Hz magnitude,  $M_{2.4}^{\rm Ma}$ . We will describe each scale in 166 167 detail in the following subsections. The MQS tracks and internally catalogs the corresponding 168 amplitudes (they are given in Table S1 in the Electronic Supplement of this article): spectral 169 plateau amplitude,  $A_0$ , for low- and high-frequency events (excluding the majority of VF events that 170 often have unclear  $A_0$  and all 2.4Hz); time domain body-wave amplitudes,  $A_P$  and  $A_S$ , at 0.2 - 0.5 171 Hz for low-frequency events (LF, BB); and 2.4 Hz spectral and time domain amplitudes,  $A_{2.4,spec}$  and 172 A<sub>2.4,pick</sub>, for high-frequency events (*HF*, *VF*, *2.4Hz*). Examples of each type are shown in Figure 2.

Unlike in Böse *et al.* (2018), we calibrate each magnitude directly against the moment magnitude,  $M_W$ , which is possible for the limited magnitude range (< 4) that we work in. We adopt the definition of moment magnitude,  $M_W$ , for Earth, which is based on the seismic moment,  $M_0$  [in MKS units] (Hanks and Kanamori, 1979):

179 
$$M_W = \frac{2}{3} (log_{10}(M_0) - 9.1)$$
(2a)

180

181 Assuming that the far-field displacement,  $u(x, \tau)$ , due to the direct S-wave from the source to the 182 receiver is a unipolar pulse and that  $M_0$  is proportional to the area under this pulse, it follows (e.g. 183 Aki & Richards, 1980)

184

185 
$$M_0 = \frac{4\pi\rho v_S^3}{R_S^{\theta\phi} s_f} R \int_0^T \boldsymbol{u}(\boldsymbol{x},\tau) d\tau = c_0 R \Omega_0$$
(2b)

186

187 with average radiation pattern  $R_s^{\theta\phi}$ , free-surface amplification  $S_b$  hypocentral distance R, density 188  $\rho$ , shear-wave velocity  $v_s$  at the source, and spectral displacement amplitude-density  $\Omega_0 = \int_0^T \boldsymbol{u}(\boldsymbol{x}, \tau) d\tau$  at frequency 0. Inserting equation (2b) in (2a) gives

190

191 
$$M_w = \frac{2}{3} (\log_{10}\Omega_0 + 1.0 \log_{10}R + \log_{10}c_0 - 9.1)$$
(2c)

192

### 193 To describe the spectral magnitude on Mars, we adopt in this study a more general form

194

195 
$$M_{W,\text{spec}}^{\text{Ma}} = \frac{2}{3} (log_{10}A_0 + a_i \ log_{10}\Delta + c_i)$$
(3a)

196

197 where  $A_0 [m/\sqrt{Hz}]$  is the square-root of the spectral amplitude extrapolated from the long-period 198 plateau of the event power spectral density (PSD) displacement spectrum. As a consequence of the 199 relation between amplitude spectral density and power spectral density, it is  $\Omega_0 = A_0/\sqrt{T}$ , where 200 *T* is the length of the data window used in the computation of the PSD. We replaced in equation 201 (3a) the hypocentral *R* [*m*] by the epicentral distance  $\Delta$  [°], assuming that the seismic ray paths 202 have a curvature similar to the curvature of the Martian surface. In equation (3a),  $log_{10}(\sqrt{T})$  as well as the logarithm of the conversion factor between meters and degrees are included in the additive coefficient  $c_i$ .

205

206 In this study,  $M_{W,\text{spec}}^{Ma}$  is set as a reference for the calibration of the other magnitude scales for 207 events with measurements of multiple amplitude types (Figure 3). This is because (1) we consider 208 the  $M_{W,\text{spec}}^{\text{Ma}}$  the most reliable measurement of the seismic moment,  $M_0$ , and thus  $M_W$ , because it is 209 estimated from the extrapolated spectrum at long periods (20 - 100 s), thus mitigating the effect 210 of attenuation; (2) due to the low SNR of many events, the picking of peak amplitudes in the time 211 domain is less robust compared to the calculation of a spectrum of a long time window, considering 212 the duration of these events generally range from 5 - 30 minutes (Figure 2); and (3) aside from the 213 majority of VF and all 2.4Hz resonance events, M<sub>W,spec</sub> can be determined for all low- and high-214 frequency events in the InSight MOS catalog (Table 1 & Figure 3).

215

216 In order to derive magnitude scales that approximate  $M_W$  and that can be applied to quantify the 217 source sizes of all events in the InSight catalog regardless of their frequency content, we will start 218 this work with an absolute calibration of  $M_{W,\text{spec}}^{\text{Ma}}$  (strictly speaking of  $M_{W,\text{spec},LF}^{\text{Ma}}$  for low-219 frequency and of  $M_{W,\text{spec},HF}^{Ma}$  for high-frequency events) relative to  $M_W$ , based on synthetic data and 220 theoretical considerations, respectively. Then, in the next step, we use the low-frequency events (LF, BB) in the InSight catalog to calibrate body-wave magnitudes,  $m_{\rm b}^{\rm Ma}$  and  $m_{\rm bs}^{\rm Ma}$ , relative to 221  $M_{W,\text{spec},LF}^{\text{Ma}}$ , and the high-frequency events (*HF*) to calibrate the 2.4 Hz magnitude,  $M_{2.4}^{\text{Ma}}$ , relative to 222  $M_{W,\text{spec},HF}^{\text{Ma}}$  (Figure 3). The resulting  $M_{2.4}^{\text{Ma}}$  scale can be applied also to the VF and 2.4Hz events to 223 224 obtain magnitude estimates that are comparable to the other events in the InSight catalog, and 225 hence ensure that every marsquake with an assigned distance also has a magnitude.

226

Our magnitude calibration is done by regressing the InSight data to determine the coefficients  $a_i$ and  $c_i$  in equation (1). This equation predicts a power decay of seismic amplitudes, i.e. a constant  $a_i$  factor for all magnitudes and distances. This assumption, however, is valid only for events outside the low (S-wave) amplitude-zone between 35° and 60° (e.g. Giardini *et al.*, 2020). Therefore, we use only S-wave amplitudes for events outside this zone for magnitude calibration 232 of  $m_{bS}^{Ma}$ .

### 234 Absolute Calibration of Spectral Magnitude M<sup>Ma</sup><sub>W,spec</sub>

We start with an absolute calibration of spectral magnitude  $M_{W,\text{spec}}^{\text{Ma}}$  (strictly speaking of  $M_{W,\text{spec},LF}^{\text{Ma}}$ 235 236 and  $M_{W,\text{spec},HF}^{Ma}$  with  $M_W$  adopting the functional form in equation (3a). The spectral magnitude can be computed for all low- and high-frequency events, excluding VFs that have unclear  $A_0$  and all 237 238 2.4Hz events, in the InSight event catalog (Figure 3) using their spectral amplitude,  $A_0$  (Figure 2). 239 The corresponding spectra are computed from the instrument-corrected displacement waveforms 240 that have been rotated into ZNE (up-down, north-south, east-west) components, from which the 241 linear trend has been removed. We apply Welch's power spectral density (PSD) method to 25.6 s 242 long Hanning-windowed time segments with 50% overlap and with zero-padding to 4096 samples 243 (102.4 s), and then compute the square-root of the resulting spectrum. From the amplitude and 244 slope between 0.1 and 0.8 Hz, the low-frequency plateau of the spectrum is extrapolated.

245

To determine the spectral amplitude,  $A_0 [m/\sqrt{Hz}]$ , we fit the following function to the spectra (vertical component)

248

$$A(f) = A_0 A_{\rm src}(f) A_{\rm att}(f)$$
(4a)

250

which combines a simple source spectrum (Brune, 1970)

252

253 
$$A_{\rm src}(f) = \frac{1}{1 + (f/f_c)^2}$$
(4b)

254

### 255 with corner frequency $f_c$ with an exponential attenuation term described by

$$A_{\text{att}}(f) = \exp(-\pi f t^*) \tag{4c}$$

259 where  $t^*$  describes the energy loss at high frequencies due to attenuation. Assuming that the corner frequency is related to the stress drop,  $f_c = 0.49 v_s (\Delta \sigma / M_0)^{1/3}$ , a shear-wave velocity of  $v_s$ 260 261 = 3.5 km/s in the source region, and a stress drop of  $\Delta \sigma$  = 1.0 MPa, the predicted corner frequency 262 of the source-time function is above 1 Hz and therefore the observed spectrum between 0.1 and 1 263 Hz is dominated by attenuation effects. We therefore fix  $f_c$  at a value of 1 Hz, so that  $A_0$  and  $t^*$  can 264 be inferred by linear inversion. For the high-frequency events, in contrast, we invert for all three 265 parameters  $(A_0, t^*, \text{ and } f_c)$ .  $t^*$  ranges from 0.6 - 4.3 s for the observed low-frequency and from 0.05 266 - 0.5 s for the observed high-frequency events (see Table S1 in Electronic Supplements). Detailed 267 subsurface structural models below the InSight lander are currently being developed (e.g. 268 Lognonné et al., 2020) and will help to fine-tune these parameters in the future (see Conclusions 269 & Outlook).

270

### 271 Spectral Magnitude $M_{W,spec,LF}^{Ma}$ for Low-Frequency Events

272 With the current interpretation that seismic waves of marsquakes in the low-frequency event 273 family have travelled through the Martian lithosphere and mantle (see Introduction; Figure 1), we 274 assume a volumetric seismic wave propagation for these events. By the conservation of energy, 275 standard geometrical spreading of seismic wave amplitudes predicts that the energy in a unit area 276 of the growing wavefront decreases as  $\Delta^2$ . Body-wave amplitudes hence decay with  $\Delta^{-1}$ , so we 277 expect  $\log_{10}(A_0)$  to decrease by about one unit per unit change in  $\log_{10}(\Delta)$ , corresponding to  $a_i \approx$ 278 1.0 (see equation 2c). Therefore, we model the decay of seismic wave amplitudes with distance 279 for the low-frequency events equation (3a) as  $a_i = (1.0 \pm 0.1)$ , assuming an uncertainty of 0.1 due 280 to leaking energy and effects of the velocity structure on geometrical spreading.

281

In order to determine the constant  $c_i$  in equation (3a) for the low-frequency events, we model a set of several hundreds of synthetic marsquakes at  $\Delta = 5^{\circ} - 150^{\circ}$  distance and down to 50 km depth for different  $M_0$  and source mechanisms using Instase (van Driel *et al.*, 2015; Ceylan *et al.*, 2017) for a set of representative Mars models described in Clinton *et al.* (2017) for the timewindow before surface-wave arrival. The same set of models and seismic propagation code was used in our original pre-launch magnitude paper (Böse *et al.*, 2018). From these simulations we determine  $c_i = (12.6 \pm 0.5)$ , and thus we obtain for the spectral magnitude for low-frequency events (*LF*, *BB*)

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291 
$$M_{W,\text{spec},LF}^{\text{Ma}} = \frac{2}{3} \left( log_{10}(A_0) + (1.0 \pm 0.1) log_{10}(\Delta) + (12.6 \pm 0.5) \right)$$
(5a)

292

The coefficients in equation (5a) differ from those in Böse *et al.* (2018), since here  $a_i$  has been set to 1.0.

295

The spectral fitting procedure to determine  $A_0$  (equation (4a)) and thus  $M_{W,\text{spec},LF}^{Ma}$  (equation (5a)) is affected by uncertainties, which can be used to estimate magnitude errors  $\sigma_{M_{W,\text{spec},LF}}^{2}$ . It follows from error propagation

299

300 
$$\sigma_{M}^{2} = \sigma_{M_{W,\text{spec},LF}}^{2} = \left(\frac{2}{3}\right)^{2} \left[ \left| \frac{\partial M_{W,\text{spec},LF}^{\text{Ma}}}{\partial \log_{10}(A_{0})} \right|^{2} \sigma_{\log_{10}(A_{0})}^{2} \right]$$

301 
$$+ \log_{10}(\Delta)^2 \left| \frac{\partial M_{W, \text{spec}, LF}^{\text{Ma}}}{\partial a_i} \right|^2 \sigma_{a_i}^2 + a_i^2 \left| \frac{\partial M_{W, \text{spec}, LF}^{\text{Ma}}}{\partial \log_{10}(\Delta)} \right|^2 \sigma_{\log_{10}(\Delta)}^2$$

$$302 \qquad \qquad + \left|\frac{\partial M_{W,\text{spec},LF}^{\text{Ma}}}{\partial c_i}\right|^2 \sigma_{c_i}^2$$

303 
$$= \frac{4}{9} \left( \sigma_{log_{10}(A_0)}^2 + 0.1 \log_{10}(\Delta)^2 + \sigma_{log_{10}(\Delta)}^2 + 0.3 \right)$$

304 = 
$$0.44 \sigma_{log_{10}(A_0)}^2 + 0.044 \log_{10}(\Delta)^2 + 0.44 \sigma_{log_{10}(\Delta)}^2 + 0.13$$
 (5b)

305

We determine  $\sigma_{log_{10}(A_0)}$  from the uncertainty (standard deviation) in spectral fitting for each lowfrequency event individually. The second term, 0.044  $log_{10}(\Delta)^2$ , suggests a distance-dependent error, ranging between 0.074 at  $\Delta = 20^\circ$  and 0.17 at  $\Delta = 90^\circ$  and that is 0.1 at  $\Delta = 32^\circ$ .  $\sigma_{log_{10}(\Delta)}$  is more difficult to determine. We use distance uncertainties as provided in the InSight catalog v.5 whenever available. For events without uncertainty estimates we assume a distance error of 25% (see Discussion). The resulting magnitude uncertainties are provided in Table S1.

### 313

### 314 Spectral Magnitude M<sup>Ma</sup><sub>W,spec,HF</sub> for High-Frequency Events

315 With the current interpretation that seismic waves of high-frequency events are guided waves 316 travelling in the crust or in a shallow layer (see Introduction; Figure 1), we assume a surface-wave 317 propagation for these events. The attenuation of surface-waves is generally more complex than 318 that of body-waves. Based on values determined by Evernden et al. (1971), Basham (1971), and 319 Nuttli (1973) for short-period Lg and Rg phases on Earth (see also Bormann et al., 2013), we set 320 in equation (3a)  $a_i = (0.8 \pm 0.1)$  for the high-frequency events. Assuming that the spectral 321 amplitudes at close distance to the seismic source (here: at  $\Delta = 1^{\circ}$ ) are similar for both low- and 322 high-frequency events, we determine from equation (5a)  $c_i = 12.8$ , and thus obtain for the 323 spectral magnitudes for high-frequency events (*HF* and *VF*)

324

325 
$$M_{W,\text{spec},HF}^{\text{Ma}} = \frac{2}{3} \left( log_{10}(A_0) + (0.8 \pm 0.1) log_{10}(\Delta) + 12.8 \right)$$
(5c)

326

327 For simplicity, we set here  $\sigma_{M_{w,spec,HF}^{Ma}} = 0.2$  for all high-frequency events, pointing out, however, 328 that the true uncertainties, in particular considering the uncertainty in the distance estimate, could 329 be much larger.

330

### 331 Calibration of Body-wave Magnitudes $m_b^{Ma}$ and $m_{bS}^{Ma}$ Relative to $M_{W,spec}^{Ma}$

In the next step, we calibrate P- and S-body-wave magnitudes,  $m_b^{Ma}$  and  $m_{bS}^{Ma}$ . These magnitudes can be determined for low-frequency events (*LF*, *BB*), excluding for  $m_{bS}^{Ma}$  the events S0325a, S0357a, S0183a, and S0205a that fall into the low (S-wave) amplitude-zone (Giardini *et al.*, 2020) (Table 1 & Figure 3). We compute  $m_b^{Ma}$  and  $m_{bS}^{Ma}$  from the peak displacement amplitudes  $A_P$  and  $A_S$ , which we determine from the integrated and filtered velocity waveforms (applying a 6<sup>th</sup> order causal Butterworth bandpass with corner frequencies of 0.167 Hz and 0.5 Hz) in the time window of the compression- or shear-wave arrivals, respectively (Figure 2b). We use the vertical component of 339 the seismograms to compute  $A_P$  and the maximum of the two horizontal components to compute 340  $A_S$ . 341 From the SNR-weighted amplitude regression of  $A_P$  and  $A_S$  with  $M_{W,\text{spec},LF}^{Ma}$ , we find for the P-wave 342 343 magnitudes 344  $m_{\rm b}^{\rm Ma} = log_{10}(A_P) + 0.73 log_{10}(\Delta) + 11.8$ 345 (6) 346 347 and for S-wave magnitudes 348  $m_{\rm bS}^{\rm Ma} = log_{10}(A_{\rm S}) + 1.06 log_{10}(\Delta) + 10.9$ 349 (7) 350 351 The MQS assigns SNRs to each (time-domain) picked amplitude (Clinton et al., 2020). Only the lowfrequency event family (LF and BB) have significant energy between 0.2 and 0.5 Hz, so  $m_{
m b}^{
m Ma}$  and 352 353  $m_{bS}^{Ma}$  can be determined only for these events (Table 1). According to the MQS v.5 event catalog < 354 20% of these events have a SNR of < 2.0 in the 0.2 - 0.5 Hz frequency band. These are mostly quality 355 'C' and 'D' events and their body-wave magnitudes are thus considered less certain.

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# 358 Calibration of 2.4Hz Magnitude $M_{2.4}^{Ma}$ Relative to $M_{W,spec}^{Ma}$

Finally, we calibrate the 2.4Hz magnitude,  $M_{2.4}^{Ma}$ , which can be computed for all high-frequency events (*HF*, *VF*, 2.4Hz) in the InSight v.5 catalog (Table 1). We assume a constant, magnitudeindependent amplification of the 2.4 Hz resonance. We calibrate the resulting magnitudes to scale with  $M_{W,spec}^{Ma}$  for all high-frequency events that have additional measurements of  $A_0$  (*HF* and *8 VF* events: S0128a, S0264e, S0334a, S0376a, S0387a, S0421a, S0424c & S0475a, but no 2.4Hz, see Table 1).

366 We define and calibrate the  $M_{2.4}^{Ma}$  in two ways:  $M_{2.4,pick}^{Ma}$  is calculated from the time domain amplitude  $A_{2.4,\text{pick}}$ ;  $M_{2.4,\text{spec}}^{\text{Ma}}$  is calculated from the spectral amplitude  $A_{2.4,\text{spec}}$ . The first is 367 368 computed from the peak amplitude on the vertical component in the 2-3 Hz bandpass-filtered (6<sup>th</sup> 369 order Butterworth bandpass with corner frequencies of 2.2 Hz and 2.8 Hz) velocity waveform; the 370 second from fitting a Lorentz curve to the displacement spectra (computed as described in the 371 spectral magnitude section) between 2 and 3 Hz (Figure 2a). 372 373 From SNR-weighted amplitude regression of  $A_{2.4,pick}$  and  $A_{2.4,spec}$  with  $M_{W,spec,HF}^{Ma}$ , we find 374 375  $M_{24\,\text{nick}}^{\text{Ma}} = log_{10}(A_{24\,\text{nick}}) + 1.0 \, log_{10}(\Delta) + 10.8$ (8a) 376 and  $M_{2.4,\text{spec}}^{\text{Ma}} = log_{10}(A_{2.4,\text{spec}}) + 1.0 \ log_{10}(\Delta) + 11.0$ 377 (8b) 378 379

### 380

### 381 **RESULTS**

We use the event distances from the InSight MQS catalog v.5 (until 2020/10/12) and amplitudes at different frequencies (as defined above) determined by MQS (Clinton *et al.*, 2020) to determine magnitudes and uncertainties using the equations described above. Table S1 (Electronic Supplement) summarizes all parameters and resulting magnitudes.

386

387 Overall, we achieve a very good agreement across the magnitude scales with standard deviations 388 of  $\sigma = 0.3$  for  $m_b^{Ma}$ ,  $m_{bs}^{Ma}$ , and of  $\sigma = 0.2$  for  $M_{2.4,\text{pick}}^{Ma}$  and  $M_{2.4,\text{spec}}^{Ma}$  with regards to the spectral 389 magnitude,  $M_{W,\text{spec}}^{Ma}$ , which we used as a calibration reference (Figure 4). There is basically no 390 difference between the sub-categories of the low- and high-frequency event families. The time 391 domain and spectral domain resonance magnitudes,  $M_{2.4,\text{pick}}^{Ma}$  and  $M_{2.4,\text{spec}}^{Ma}$ , are very similar, while 392  $M_{2.4,\text{spec}}^{Ma}$  achieves a slightly better fit with  $M_{W,\text{spec}}^{Ma}$ . We conclude that the two spectral magnitudes, 393  $M_{W,\text{spec}}^{Ma}$  and  $M_{2.4,\text{spec}}^{Ma}$ , are most robust and we thus set the preferred (moment) magnitude to 394  $M_W^{Ma} = M_{W,spec,LF}^{Ma}$  for *LF* and *BB* events, to  $M_W^{Ma} = M_{W,spec,HF}^{Ma}$  for *HF* events, and to  $M_W^{Ma} = M_{2.4,spec}^{Ma}$ 395 for *VF* and 2.4Hz resonance events. The picking of peak amplitudes in the frequency domain is 396 more stable, since a much longer time-window (5 - 30 minutes) is considered (Figure 2).

397

398 Adopting the newly calibrated scales (Table 2), we determine the magnitudes of the events in the 399 InSight MQS catalog v.5 as  $1.1 \le M_W^{Ma}(\sim M_W) \le 3.7$  (Table S1). More specifically, we find  $2.7 \le M_W^{Ma} \le 3.7$  for the 20 *LF* events,  $2.7 \le M_W^{Ma} \le 3.5$  for the 9 *BB* events,  $1.7 \le M_W^{Ma} \le 2.6$  for 400 the 50 HF events,  $1.1 \le M_W^{Ma} \le 2.9$  for the 21 VF events, and  $1.3 \le M_W^{Ma} \le 3.0$  for the 179 2.4Hz 401 402 events. From equation (5b) we estimate the magnitude uncertainties for the low-frequency (LF, 403 *BB*) events to range between  $\sigma_M = 0.2$  and  $\sigma_M = 0.5$  (Table S1). As stated before, because of the 404 currently unclear interpretation of high-frequency events, including their distances (Giardini et al., 405 2020; van Driel *et al.*, 2021), we simply set  $\sigma_M = 0.2$ , even though the true uncertainties could be 406 much larger. Magnitude estimates for high-frequency events given in this article primarily serve 407 for a relative comparison of these events with each other (see Discussion).

408

Figure 5 shows the distribution of all events as a function of epicentral distance. The four largest events are all low-frequency events: S0325a ( $M_W^{Ma}3.7$ ,  $\Delta = 38^\circ$ , quality 'B'), S0167a ( $M_W^{Ma}3.7$ ,  $\Delta =$ 95°, quality 'C'), S0173a ( $M_W^{Ma}3.7$ ,  $\Delta = 28^\circ$ , quality 'A'), and S0235b ( $M_W^{Ma}3.5$ ,  $\Delta = 27^\circ$ , quality 'A'). Note that the distance estimate for S0167a is very uncertain, since it was derived from waveform envelope similarity (see Giardini *et al.*, 2020 for details). The large uncertainty in distance for S0167a is also reflected in inconsistent magnitude estimates across different scales ( $m_{bS}^{Ma}3.3$  and  $M_{W,spec}^{Ma}3.7$ , see Table S1).

416

417

#### 418 **DISCUSSION**

In this article, we recalibrated magnitude relations for Mars from InSight data using a seismicity catalog of a full Martian year, corresponding to nearly two Earth years. Pre-landing seismicity estimates expected larger marsquakes that more closely resembled teleseismic events we observe on Earth. Hence, our pre-landing magnitude study assumed larger events, with surface-waves and 423 energy at long periods (Böse *et al.*, 2018). An additional crucial difference from our pre-landing
424 study is that the available data has included the routine collection of continuous high-sampled rate
425 data (first 10 sps, then 20 sps) throughout the mission, instead of the 2 sps planned beforehand.
426 This means higher frequency data than anticipated could be used. The high-frequency event class
427 had not been expected at all and therefore required completely novel magnitude scales, as
428 presented in this article.

429

(Single-station) magnitude computations are affected by assumptions made to model the seismic wave propagation from the source to the receiver. These models predict both the decay of seismic wave amplitudes and the observation of different phase arrivals at different distances, and they also impact the distance estimates relative to the receiver (Khan *et al.*, 2016). Despite collecting an exceptional seismic event dataset during the first year of the InSight mission (InSight Mars SEIS Data Service, 2019), our knowledge of Mars interior structure is still limited, which leads to unavoidable uncertainties in magnitude estimates.

437

438 Model assumptions have a strong impact on magnitude estimates as shown for the Moon. 439 Nakamura et al. (1976) studied the distant-dependent decay of S-wave amplitudes on the Moon 440 generated by artificial and meteoroid impacts as well as shallow moonquakes observed at the 441 Apollo seismic network stations. The authors assigned lunar magnitudes to these events, 442 estimated to be approximately 1.0 less than the Richter magnitude for the same amplitude. Goins 443 et al. (1981) developed equations to estimate source parameters of shallow and deep moonquakes 444 (seismic moment, seismic energy release, stress drop, and body-wave magnitude), from 445 seismograms and displacement spectra that account for the effects of instrument bandwidth, 446 variations in corner frequency, and the effects of intense scattering on the Moon. They determined 447 the stress drops and body-wave magnitudes of the largest shallow moonquakes as 40 MPa, and  $\sim$ 448 5.0, respectively. Following the same method with some improvements and corrections, Oberst 449 (1987) obtained remarkably different results with stress drops of > 100 MPa and body-wave 450 magnitudes of > 5.5.

451

In our current understanding (Giardini *et al.*, 2020; van Driel *et al.*, 2021) the seismic waveforms of events in the low-frequency (*LF*, *BB*) and high-frequency (*HF*, *VF*, *2.4Hz*) event families are dominated by mantle and crustal guided waves, respectively (Figures 1 and 2). Following this interpretation, we constrained the magnitude scales in this article, assuming a volumetric seismic wave propagation for the low-frequency and a surface propagation for the high-frequency events.

We determined the constant in equation (5a) from seismic waveform modeling through a set of representative Mars models discussed in Clinton *et al.* (2017). From equations (2b) and (2c) and  $\Omega_0 \approx A_0 \sqrt{N\Delta t}$ , where *N* is the number of data samples before zero-padding in the time windows and  $\Delta t$  the sampling interval of the displacement time series used for the PSD computation, we would theoretically expect for S-waves that

463

$$c_i = \log_{10}(\sqrt{N\Delta t}) + a_i \log_{10}\left(\frac{\pi}{180}R_{Ma}\right) + \log_{10}\frac{4\pi\rho v_S^3}{R_S^{\theta\phi}S_f} - 9.1$$
(9)

465

464

466 With *N* = 512,  $\Delta t$  = 0.05 s,  $a_i$  = 1.0, density  $\rho$  = 3,500  $kg/m^3$ ,  $v_s$  = (3,500 ± 500) m/s, average (Swave) radiation pattern of  $R_S^{\theta\phi} = 0.55$ , free-surface amplification  $S_f = 2.0$  (Boore & Boatwright, 467 468 1984), and radius of Mars  $R_{Ma}$  = 3'385 km, we obtain  $c_i$  = 0.7 + 4.8 + (15.2 ± 0.2) - 9.1 = (11.6 469  $\pm$  0.2). The discrepancy between our modelled ( $c_i = 12.6$ ) and theoretical value ( $c_i = 11.6 \pm 0.2$ ) 470 suggests that the theoretical approximation in equation (2b) is likely too simple for our purpose, 471 since we are not looking at a single seismic ray from the source through a full space, but at a time 472 window that contains both surface reflections and interfering rays. In fact, at the epicentral 473 distances of the observed marsquakes, the seismograms can no longer be regarded as the 474 deterministic signal that is assumed by the theory underlying equation (2b), but must be treated 475 as transient stochastic time series with variable amplitudes and undetermined duration. 476 Additional simulations based on the convolution of a synthetic source-time function of known 477 seismic moment with a transient white-noise time series multiplied with a tapering function that 478 decays exponentially with time show that the spectral analysis over a more or less arbitrary length 479 of the S-wave coda tends to underestimate the true seismic moment. Thus, at least part of the 480 difference between the theoretically expected value of  $c_i$  and the value of  $c_i$  derived from the

481 calibration procedure based on synthetic seismograms, implemented in the present study,482 compensates for this methodological bias.

483

484 Despite the apparent disagreement with the theoretically expected  $c_i$  coefficient, the adequacy of 485 our magnitude scales is supported two-fold: (1) In an independent study, Brinkman et al. (2020) 486 determined the magnitudes of the S0235b, S0173a, and S0183a events from moment tensor 487 inversion as  $M_W$  3.1,  $M_W$  3.0, and  $M_W$  3.3, respectively, which is close to our values ( $M_W^{Ma}$  3.5 ± 0.3, 488  $M_W^{Ma}3.7 \pm 0.4$ , and  $M_W^{Ma}3.1 \pm 0.4$  respectively; see Table S1). The remaining discrepancy can be 489 explained, since for S0235b and S0173a, the InSight station is located close to a maximum of the 490 S-wave radiation pattern (as estimated from the amplitude 491 ratios of P- and S-waves on different components), resulting in an overestimation of magnitude 492 when not inverting for the moment tensor. (2) Figure 6 compares our  $m_b^{Ma}$  and  $m_{bS}^{Ma}$  estimates 493 from  $A_P$  and  $A_S$  for the low-frequency (LF, BB) events in the InSight catalog with the medium-494 period body-wave broadband magnitude, mB, proposed by Gutenberg (1945a,b) for Earth. 495 Although developed for teleseismic P- and S-waves at 2 - 20 s period on Earth (Bormann et al., 496 2013), we are applying the *mB* relations here to Mars for  $A_P$  and  $A_S$  at 2 - 6 s period. Simplifying 497 the lookup table in Gutenberg and Richter (1956) for a source depth of 15 km and two distance 498 intervals, we obtain

499

500 $mB_P = \log_{10}(A_P) + 12.9 \text{ for } \Delta \le 70^o$	(10a)
---	-------

$mB_P = log_{10}(A_P) + 13.2$ for $\Delta > 70^o$	(10b)
	$mB_P = log_{10}(A_P) + 13.2$ for $\Delta > 70^o$

502 
$$mB_S = log_{10}(A_S) + 12.6 \text{ for } \Delta \le 70^o$$
 (10c)

 $mB_{\rm S} = log_{10}(A_{\rm S}) + 13.0$  for  $\Delta > 70^{\circ}$ 

503

504

The good agreement of  $m_b^{Ma}$  and  $m_{bs}^{Ma}$  with  $mB_p$  and  $mB_s$  in Figure 6 suggests that the attenuation of body-waves at long periods and teleseismic distances on Mars and Earth are similar. While we could potentially consider applying earthquake magnitude scales, such as mB, to Mars, it needs to be pointed out again that due to the noise condition on Mars, time-domain amplitudes  $A_p$  and  $A_s$ could so far be determined only for a small subset of marsquakes. For the majority of seismic events we have only measurements of spectral amplitudes ( $A_0$ ,  $A_{2.4,spec}$ ) for which no comparable

(10d)

511 magnitude scales on Earth exist. On the other hand, the good agreement of mars- and earthquake 512 magnitude scales in Figure 6 indicates that both the absolute calibration of  $M_{W,spec}^{Ma}$  and the 513 relative calibration of (at least) of  $m_{bs}^{Ma}$  and  $m_{bs}^{Ma}$  achieve reasonable magnitude estimates.

514

515 All marsquake signals identified by MQS so far have durations of 5 to 30 minutes. This is 516 significantly longer compared to what we would expect on Earth for the same magnitude range (M 517 < 4), even though the comparison is difficult since such small events are usually not observed 518 in distances of > 1500 km. We attribute these long durations mainly to the strong scattering of 519 seismic energy on Mars (e.g., Lognonné et al., 2020). On Earth, the signal length has been used to 520 compute duration (or coda) magnitudes for strong historical and local earthquakes (Bormann et 521 al., 2013). On Mars, such a relationship, however, is not clearly visible, likely because scattering is 522 the dominating effect on the signal duration, which is independent of the source magnitude.

523

524 Distances of low-frequency (LF, BB) events relative to the InSight lander in the event catalog 525 (InSight Marsquake Service, 2021) were determined using the probabilistic approach of Böse et al. 526 (2016), which accounts for both pick and model uncertainties. This approach requires at least 2 527 seismic phase picks. Absolute event locations in the catalog (including a backazimuth) were 528 inferred only for events with a clear polarization of the seismic signals, which only applies to a 529 small subset of events (Clinton et al., 2020). As discussed in Böse et al. (2018), even large distance 530 uncertainties on the order of 25% and depth errors of 20 km (L1 InSight mission requirements; 531 Banerdt *et al.*, 2013) translate into magnitude errors of only  $\pm$  0.2 units, which is smaller than the 532 magnitude uncertainties ( $\pm$  0.5 units) that can generally be expected with a single-station approach 533 and that are related to the variability in the focal mechanisms and uncertainties in the interior 534 model.

535

A number of low-frequency events without clear emergent phase arrivals in the InSight catalog were located based on the similarity of waveform envelopes and theoretical travel-time curves using locatable events with clear phases as a reference (Giardini *et al.*, 2020). Estimating the distance (and thus magnitude) uncertainties for these events is difficult. They might be quite large, in particular for events at large distances. Phase identification for these events is ambiguous, some

541 of them have extremely weak P- and S-wave signals (Giardini et al., 2020). For instance, it is 542 possible that  $A_P$  (and potentially  $A_S$ ) for the most distant marsquake (S0167a) at 95° distance was 543 measured in a different phase than for the closer events, or that the event is even beyond the core 544 shadow and we observe a SS/SSS wave train. This would also explain the different shape of the 545 waveform envelopes as observed in Giardini et al. (2020). For this reason, we consider the current 546 magnitude estimates for S0167a and other events in this distance range to be rather uncertain. 547 Further research is needed to better understand the seismic wave propagation through Mars, 548 which will lead to a better phase assignment and enhanced distance and thus magnitude estimates 549 in the future.

550

551 Locating the (crustal) high-frequency (HF, VF, 2.4Hz) events requires another approach. In the 552 InSight catalog distances of these events were inferred from P<sub>g</sub>-S<sub>g</sub> travel-times using a very simple 553 model with constant velocities of  $v_p = 4$  km/s and  $v_s = v_P/\sqrt{3}$  (= 2.3 km/s) (van Driel *et al.*, 554 2021). With this model, high-frequency events are located at distances between 3° and 40°, with 555 the bulk between 25 and 30°, while low-frequency events are located at 20° - 95° distance (Figure 556 5). Assuming lower (e.g.  $v_p = 2.6$  km/s and  $v_s = 1.5$  km/s) or higher (e.g.  $v_p = 5.2$  km/s and  $v_s = 1.5$  km/s) or higher (e.g.  $v_p = 5.2$  km/s and  $v_s = 1.5$  km/s) or higher (e.g.  $v_p = 5.2$  km/s and  $v_s = 1.5$  km/s) or higher (e.g.  $v_p = 5.2$  km/s and  $v_s = 1.5$  km/s) or higher (e.g.  $v_p = 5.2$  km/s and  $v_s = 1.5$  km/s) or higher (e.g.  $v_p = 5.2$  km/s and  $v_s = 1.5$  km/s) or higher (e.g.  $v_p = 5.2$  km/s and  $v_s = 1.5$  km/s) or higher (e.g.  $v_p = 5.2$  km/s and  $v_s = 1.5$  km/s) or higher (e.g.  $v_p = 5.2$  km/s) or higher ( 557 3.0 km/s) velocities, the distances of the high-frequency events would change by approximately 558 10° (1° on Mars corresponds to 59.2 km). Using equation (6b), this translates again into magnitude 559 uncertainties of  $\pm$  0.2 units, similarly to the expected uncertainties for low-frequency events, and 560 again smaller compared to what can generally be expected with a single station approach ( $\pm 0.5$ 561 units). Because of the currently unclear interpretation of events in the high-frequency family, 562 including large uncertainties in their distances, our magnitude estimates for these events serve 563 primarily to enable a relative comparison of these events with each other.

564

Models of the subsurface structure of Mars below the InSight lander start emerging. For example, the receiver function analysis of Lognonné *et al.* (2020) suggests a 10 km low-velocity layer, which might generate some type of site effect. The 2.4 Hz resonance is probably caused by even shallower layers, but these are most likely not the ones in which the guided waves of the high-frequency events propagate (van Driel *et al.*, 2021). Nevertheless, the propagation velocity of the guided waves could significantly differ from those we assumed at the source, resulting in a systematic bias in the moment estimation. Incorporating site effects as predicted from this and other models, which are currently being developed, could help to fine-tune magnitude estimates on Mars in the future.

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### 576 CONCLUSIONS & OUTLOOK

577 In this article, we revised the magnitude relations in Böse et al. (2018) to account for the seismic 578 and noise characteristics observed on Mars during the first Martian year of the InSight mission. In 579 light of these observations, we updated the scaling relations for the spectral and body-wave magnitudes,  $M_{W,spec}^{Ma}$ ,  $m_b^{Ma}$  and  $m_{bs}^{Ma}$ , and introduced a new magnitude scale,  $M_{2,4}^{Ma}$ , for high-580 581 frequency events. The new scales are summarized in Table 2 of this article. We recommend setting the preferred magnitude to  $M_W^{\text{Ma}} = M_{W,\text{spec},LF}^{\text{Ma}}$  for LF and BB events, to  $M_W^{\text{Ma}} = M_{W,\text{spec},HF}^{\text{Ma}}$  for HF 582 events, and to  $M_W^{Ma} = M_{2.4,spec}^{Ma}$  for VF and 2.4Hz resonance events. We used the new magnitude 583 584 scales to determine that the magnitudes of the events in the current InSight catalog v.5 range 585 between 1.1 and 3.7. We estimate that magnitude uncertainties are on the order of  $\pm$  0.5 magnitude 586 units, which is mainly due to uncertainties in the focal mechanisms rather than in distances. 587 Location uncertainties of some events in the InSight catalog, however, might be quite large, in 588 particular for events at large distance with uncertain phase identification. Estimating the 589 magnitude uncertainties of these events is difficult, in particular because other phase arrivals 590 might be measured than what was used for magnitude calibration. Overall, we achieve a very good 591 agreement across all magnitude scales with regards to the spectral magnitude,  $M_{w,spec}^{Ma}$ , which is 592 considered a measure of the moment magnitude,  $M_W$ .

593

Magnitude estimates of marsquakes are a key component in the characterization of the seismic activity rate of Mars (Knapmeyer *et al.*, 2019). All pre-mission Martian seismicity models are based on estimated moment rates, which require magnitudes. Hence, magnitudes are needed to determine which models are best suited for Mars. Refining interior models of Mars from InSight observations will not only help to further refine magnitude scales for marsquakes in the future,

- 599 but also to better understand the level and distribution of seismicity on the planet. Once we have
- 600 a better understanding of the source and propagation, in particular of VF and SF events, new
- 601 magnitudes scales for these events will be devised.
- 602
- 603

### 604 DATA & RESOURCES

- 605 The Electronic Supplement to this article contains source parameters of marsquakes as provided
- in the InSight catalog v.5 (until 2020/10/12) of the InSight Marsquake Service (2021), as well as
- 607 time-domain and spectral amplitudes, spectral fitting parameters and magnitudes as computed in
- this article.
- 609 InSight Marsquake Service (2021). Mars Seismic Catalogue, InSight Mission; V5 2021-01-04. ETHZ,
- 610 IPGP, JPL, ICL, MPS, Univ. Bristol. <u>https://doi.org/10.12686/a10.</u>
- 611 InSight Mars SEIS Data Service (2019). SEIS raw data, Insight Mission. IPGP, JPL, CNES, ETHZ, ICL,
- 612 MPS, ISAE-Supaero, LPG, MFSC. <u>https://doi.org/10.18715/SEIS.INSIGHT.XB 2016.</u>
- 613 InSight Mars SEIS Data Service (2019). InSight SEIS Data Bundle. PDS Geosciences (GEO) Node.
- 614 <u>https://doi.org/10.17189/1517570</u>.
- 615 Please be aware that the InSight catalog can change over time reflecting our respective
- 616 understanding of Mars seismicity and its interior structure. The most recent event InSight catalog
- 617 can be retrieved from <u>https://www.iris.edu/hq/sis/insight</u>.
- 618
- 619

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### 632 **REFERENCES**

- Aki, K., & Richards, P. G. (1980). *Quantitative seismology: Theory and methods*. San Francisco: W.H.
  Freeman.
- Banerdt, W. B., S. Smrekar, P. Lognonné, T. Spohn, S. W. Asmar, D. Banfield, L. Boschi, U. Christensen,
  V. Dehant, W. Folkner, et al. (2013). InSight: A discovery mission to explore the interior of Mars,
  Lunar and Planetary Science Conference, Lunar and Planetary Institution. Technical Report,
  Vol. 44, 1915 pp.
- Banerdt, W.B., S. Smrekar, P. Lognonné, D. Giardini, et al. (2020). Early Results from the InSight
  Mission: Surface Environment and Global Seismic Activity, *Nature Geoscience*, 13, 183–
  189(2020), doi: 10.1038/s41561-020-0544-y.
- 642 Basham, P. W. (1971). A New Magnitude Formula for Short Period Continental Rayleigh 643 Waves, *Geophys. J. Int.* 23(2), 255–260, https://doi.org/10.1111/j.1365-246X.1971.tb01817.x.
- 644 Bisztricsany, E. (1958). A new method for the determination of the magnitude of earthquake. Geofiz.
  645 Kozl., 7, 69-96 (in Hungarian with English abstract).
- Brune, J. N. (1970). Tectonic stress and the spectra of seismic shear waves from earthquakes, *J. Geophys. Res.*, 75(26), 4997–5009, doi:10.1029/JB075i026p04997.
- Böse, M., Clinton, J. F., Ceylan, S., Euchner, F., van Driel, M., Khan, A., *et al.* (2016). A Probabilistic
  Framework for Single-Station Location of Seismicity on Earth and Mars. *Physics of the Earth and Planetary Interiors, 262,* 48–65. https://doi.org/10.1016/j.pepi.2016.11.003
- Böse, M., Giardini, D., Stähler, S. C., Ceylan, S., Clinton, J. F., van Driel, M., et al. (2018). Magnitude
  Scales for Marsquakes. *Bull. Seismol. Soc. Am.* 108(5A), 2764–2777.
  https://doi.org/10.1785/0120180037
- Boore, D., and J. Boatwright (1984). Average body-wave radiation coefficients, *Bull. Seism. Soc. Am.*,
   74, 1615-1621.
- Bormann, P., S. Wendt, D. DiGiacomo (2013). Seismic sources and source parameters, In: Bormann,
  P. (Ed.), New Manual of Seismological Observatory Practice 2 (NMSOP2), Potsdam, Deutsches
  GeoForschungsZentrum GFZ, 1-259, <u>https://doi.org/10.2312/GFZ.NMSOP-2 ch3</u>.
- Brinkman, N., S. Stähler, D. Giardini, C. Schmelzbach, A. Jacob, N. Fuji, C. Perrin, P. Lognonné, M.
  Böse, B. Knapmeyer-Endrun, E. Beucler, E. et al. (2020). Single-station moment tensor inversion
  on Mars, J. Geophys. Res. Planets, 10.1002/essoar.10503341.1.
- 662 Ceylan, S., van Driel, M., Euchner, F., Khan, A., Clinton, J., Krischer, L., Böse, M., Stähler, S. and D.
  663 Giardini (2017). From initial models of seismicity, structure and noise to synthetic
  664 seismograms for Mars, *Space Science Reviews*, InSight special issue, DOI 10.1007/s11214-017665 0380-6.
- 666 Ceylan, S., J. Clinton, D. Giardini, M. Böse, C. Charalambous, et al. (2020). Companion guide to the
   667 Marsquake catalog from InSight, sols 0–478: Data content and non-seismic events, *Physics of* 668 *the Earth and Planetary Interiors*, <u>https://doi.org/10.1016/j.pepi.2020.106597</u>
- 669 Charalambous, C., Stott, A. E., Pike, W. T., McClean, J. B., Warren, T., Spiga, A., et al. (2021). A
  670 Comodulation Analysis of Atmospheric Energy Injection into the Ground Motion at InSight,
  671 Mars. Journal of Geophysical Research: Planets 126,
  672 e2020JE006538. https://doi.org/10.1029/2020JE006538
- 673 Clinton, J. F., Giardini, D., Böse, M., Ceylan, S., van Driel, M., Euchner, F., *et al.* (2018). The Marsquake
  674 Service: Securing Daily Analysis of SEIS Data and Building the Martian Seismicity Catalogue for
  675 InSight. *Space Science Reviews*, *214*(8), 133. doi:10.1007/s11214-018-0567-5.
- 676 Clinton, J. F., D. Giardini, P. Lognonné, B. Banerdt, M. van Driel, M. Drilleau, N. Murdoch, M. Panning,
  677 R. Garcia, D. Mimoun, M. Golombek, J. Tromp, R. Weber, M. Böse, S. Ceylan, I. Daubar, B. Kenda,
  678 A. Khan, L. Perrin, and A. Spiga (2017). Preparing for InSight: an invitation to participate in a
  679 blind test for Martian seismicity, *Seismol. Res. Lett.* 88(5), 1290-1302, doi:
  680 10.1785/0220170094.
- 681 Clinton, J.F., S. Ceylan, M. van Driel, D. Giardini, S.C. Stähler, M. Böse *et al.* (2020). The Marsquake
  682 Catalogue from InSight, Sols 0-478, *Physics of the Earth and Planetary Interiors*, 683 <u>https://doi.org/10.1016/i.pepi.2020.106595</u>.
- Dahmen, N. L., Clinton, J. F., Ceylan, S., van Driel, M., Giardini, D., Khan, A., et al. (2021). Super
   high frequency events: A new class of events recorded by the InSight seismometers on

- 686Mars. JournalofGeophysicalResearch:Planets126,687e2020JE006599. <a href="https://doi.org/10.1029/2020JE006599">https://doi.org/10.1029/2020JE006599</a>
- Dahmen, N., G. Zenhäusern, et al. (subm.): Resonances and Lander Modes observed
  by InSight on Mars (1-9~Hz), submitted to BSSA.
- 690 Giardini, D., P. Lognonné, W. Banerdt, W. Pike, U. Christensen, S. Ceylan, J. Clinton, M. van Driel, S.
  691 Stähler, M. Böse *et al.* (2020). The Seismicity of Mars, *Nature Geoscience*, 13, 1-8, doi: 10.1038/s41561-020-0539-8.
- 693 Goins, N.R., A.M. Dainty, and M.N. Toksöz (1981). Seismic energy release of the Moon, *J. Geophys.* 694 *Res.* 86: 378–388.
- 695 Gutenberg, B. (1945a). Amplitudes of P, PP, and S and magnitude of shallow earthquakes. *Bull.*696 *Seism. Soc.* Am. 35, 57-69.
- 697 Gutenberg, B. (1945b). Magnitude determination of deep-focus earthquakes. *Bull. Seism. Soc. Am.*,
  698 35, 117-130.
- 699 Gutenberg, B., and Richter, C. F. (1956). Magnitude and energy of earthquakes, *Annali di Geofisica*, 9, 1-15.
- 701 Hanks, T.C. and H. Kanamori (1979). A Moment magnitude scale, *J. Geophys. Res.* 84 (B5), doi:10.1029/JB084iB05p02348.
- InSight Mars SEIS Data Service (2019). SEIS raw data, Insight Mission. IPGP, JPL, CNES, ETHZ, ICL,
   MPS, ISAE-Supaero, LPG, MFSC. <u>https://doi.org/10.18715/SEIS.INSIGHT.XB\_2016.</u>
- InSight Mars SEIS Data Service (2019). InSight SEIS Data Bundle. PDS Geosciences (GEO) Node.
   <u>https://doi.org/10.17189/1517570</u>.
- InSight Marsquake Service (2021). Mars Seismic Catalogue, InSight Mission; V5 2021-01-04. ETHZ,
   IPGP, JPL, ICL, MPS, Univ. Bristol. <u>https://doi.org/10.12686/a10.</u>
- Kennett, B. L. N. and T. Furumura (2013). High-frequency *Po/So* guided waves in the oceanic
  lithosphere: I—long-distance propagation, *Geophysical Journal International* 195(3), 1862–
  1877, <u>https://doi.org/10.1093/gji/ggt344</u>.
- 712 Khan, A., van Driel, M., Böse, M., Giardini, D., Ceylan, S., Yan, J., . . . Banerdt, W. B. (2016). Single-713 station and single-event marsquake location and inversion for structure using synthetic 714 Martian waveforms, Phys. Earth Planet. Inter. 258, 28-42. doi: 715 https://doi.org/10.1016/j.pepi.2016.05.017.
- Knapmeyer, M., B. Knapmeyer-Endrun, A.-C. Plesa, M. Böse, T. Kawamura, J.F. Clinton, M.P.
  Golombek, S. Kedar, S. Stähler, J. Stevanović, C. Perrin, P. Lognonné, N.A. Teanby, and R. Weber
  (2019). Estimation of the Seismic Moment Rate from an Incomplete Seismicity Catalog, in the
  Context of the InSight Mission to Mars, *Bull. Seismol. Soc. Am.*, 109 (3), 1125-1147,
  https://doi.org/10.1785/0120180258.
- Lognonné, P., W B. Banerdt, D. Giardini, W. T. Pike, U. Christensen, and et al. (2020). First constraints on the shallow elastic and anelastic structure of Mars from InSight data, *Nature Geoscience*, 13, 213-220, doi: 10.1038/s41561-020-0536-y.
- Lognonné, P., Banerdt, W. B., Giardini, D., Pike, W. T., Christensen, U., Laudet, P., et al. (2019). SEIS:
  Insight's Seismic Experiment for Internal Structure of Mars. *Space Science Reviews*, *215*(1), 12.
  https://doi.org/10.1007/s11214-018-0574-6
- Mimoun, D., Murdoch, N., Lognonné, P., Hurst, K., Pike, W.T., Hurley, J., Nebut, T. & Banerdt, W.B.
  (2017). The Noise Model of the SEIS Seismometer of the InSight Mission to Mars, *Space Science Review*, doi: 10.1007/s11214-017-0409-x.
- Nakamura, Y., Duennebier, F. K., Latham, G. V., and Dorman, H. J. (1976). Structure of the lunar
   mantle, *J. Geophys. Res.*, 81(26), 4818–4824, doi:<u>10.1029/JB081i026p04818</u>.
- Nissen-Meyer, T., van Driel, M., Stähler, S. C., Hosseini, K., Hempel, S., Auer, L., Colombi, A., &
  Fournier, A. (2014). AxiSEM: broadband 3-D seismic wavefields in axisymmetric media. Solid
  Earth, 5(1), 425–445. https://doi.org/10.5194/se-5-425-2014.
- Nuttli, O. W. (1973). Seismic wave attenuation and magnitude relations for eastern North
   America, J. Geophys. Res., 78(5), 876–885, doi:10.1029/JB078i005p00876.
- 737 Oberst, J. (1987). Unusually high stress drop associated with shallow moonquakes, *J. Geophys. Res.*738 92: 1397–1405.
- van Driel, M., Ceylan, S., Clinton, J. F., Giardini, D., Horleston, A., Margerin, L., et al. (2021). Highfrequency seismic events on Mars observed by InSight. *Journal of Geophysical Research: Planets*, 126, e2020JE006670. <u>https://doi.org/10.1029/2020JE006670</u>

- van Driel, M., Krischer, L., Stähler, S. C., Hosseini, K., and T. Nissen-Meyer (2015). Instaseis: instant
- global seismograms based on a broadband waveform database, *Solid Earth* 6, 701-717,
   <u>http://dx.doi.org/10.5194/se-6-701-2015</u>.

## 745 Mailing Address for each Author

5,	aren bose, institute of deophysics, owiss beisinological bervice (obb.), in in zarien, bonneggsti.
	8092 Zurich, Switzerland; <u>mboese@sed.ethz.ch</u> , <u>https://orcid.org/0000-0003-4639-719X</u>
Si	mon Stähler, Institute of Geophysics, ETH Zürich, Sonneggstrasse 5, 8092 Zürich, Switzerland;
<u>si</u>	mon.staehler@erdw.ethz.ch, https://orcid.org/0000-0002-0783-2489
N	icholas Deichmann, Swiss Seismological Service (SED), ETH Zurich, Sonneggstr, 5, 8092 Zurich,
S	vitzerland: deichmann@sed.ethz.ch
0	
р	omenico Giardini Institute of Geophysics FTH Zurich Sonneggstr 5, 8092 Zurich Switzerland
d	pmenico giardini@erdw.etbz.ch. https://orcid.org/0000-0002-5573-7638
<u>u</u>	<u>memee.giarumweruw.etuz.eu, meps.//oreiu.org/0000/0002/3373/7030</u>
L	hn Clinton Swice Saismalagical Sarvica (SED) ETH Zurich Sannaggetr 5, 2002 Zurich
JU C.	uit childen, Swiss Seisihological Service (SED), ETH Zurich, Sonneggsti. 5, 8092 Zurich,
31	vitzeriand; j <u>ennton@sed.etnz.en</u> , <u>nitips://oreid.org/0000-0001-8626-2705</u>
D	
<b>P</b>	<b>hippe Lognonne</b> , Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005
Pa	iris, France; <u>lognonne@ipgp.fr</u> , <u>https://orcid/org/0000-0002-1014-920X</u>
Sa	was Ceylan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland;
<u>Sa</u>	<u>vas.ceylan@erdw.ethz.ch;</u>
Μ	artin van Driel, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland;
Vä	ndriel@erdw.ethz.ch, https://orcid.org/0000-0002-8938-4615
C	onstantinos Charalambous, Department of Electrical and Electronic Engineering, Imperial
С	ollege London. South Kensington Campus. London. SW7 2AZ. United Kingdom:
<u>C(</u>	
	onstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895
	instantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895
N	<pre>isonstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</pre>
N ni	ikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747
N <u>ni</u>	ikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747
N ni A	<b>ikolaj Dahmen</b> , Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747 <b>nna Horleston</b> , School of Earth Sciences, University of Bristol, Wills Memorial Building, Oueens
N ni A	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>ikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RL UK: anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-</li> </ul>
N ni A R	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>ikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; <u>anna.horleston@bristol.ac.uk</u>, <u>https://orcid.org/0000-0002-6748-</u> 522</li> </ul>
N ni A R 6	<b>ikolaj Dahmen</b> , Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747 <b>nna Horleston</b> , School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens oad, Bristol BS8 1RJ, UK; <u>anna.horleston@bristol.ac.uk</u> , <u>https://orcid.org/0000-0002-6748-522</u>
N ni A R 6 S	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>akolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748- 522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, E-75005</li> </ul>
N ni A R 6 S	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>akolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748- 522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 uris. Erance: kawamura@ingn.fr.</li> </ul>
N ni A R <u>6</u> T P	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>ikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 uris, France; kawamura@ipgp.fr</li> </ul>
N ni A R G C T P	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>ikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>nna Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 aris, France; kawamura@ipgp.fr</li> </ul>
N ni A R G G T P C A	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>aikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 uris, France; kawamura@ipgp.fr</li> <li>mir Khan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, witzerland, alabar@ethz.ch, https://orcid.org/0000.0002.4462.2172</li> </ul>
N ni A R ( <u>6</u>	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>aikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 aris, France; kawamura@ipgp.fr</li> <li>mir Khan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, vitzerland; akhan@ethz.ch, https://orcid.org/0000-0003-4462-3173</li> </ul>
N ni A R G G S V S V	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>akolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 aris, France; kawamura@ipg.fr</li> <li>mir Khan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, vitzerland; akhan@ethz.ch, https://orcid.org/0000-0003-4462-3173</li> </ul>
N ni Ra 6. Ta Pa A Sv M	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>ikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 aris, France; kawamura@ipgp.fr</li> <li>mir Khan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, vitzerland; akhan@ethz.ch, https://orcid.org/0000-0003-4462-3173</li> <li>artin Knapmeyer, DLR Institute of Planetary Research, Rutherfordstr. 2, 12489 Berlin,</li> </ul>
N ni A R G C S V M G	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>akolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 aris, France; kawamura@ipgp.fr</li> <li>mir Khan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, vitzerland; akhan@ethz.ch, https://orcid.org/0000-0003-4462-3173</li> <li>artin Knapmeyer, DLR Institute of Planetary Research, Rutherfordstr. 2, 12489 Berlin, armany; martin.knapmeyer@dlr.de, https://orcid.org/0000-0003-0319-2514</li> </ul>
N ni R( <u>6</u> T P; N Sy G	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>akolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 aris, France; kawamura@ipgp.fr</li> <li>mir Khan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, vitzerland; akhan@ethz.ch, https://orcid.org/0000-0003-4462-3173</li> <li>artin Knapmeyer, DLR Institute of Planetary Research, Rutherfordstr. 2, 12489 Berlin, armany; martin.knapmeyer@dlr.de, https://orcid.org/0000-0003-0319-2514</li> </ul>
N ni R ( <u>6</u> T ( P ( N S V M G ( G	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>ana Horleston, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 aris, France; kawamura@ipgp.fr</li> <li>mir Khan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, vitzerland; akhan@ethz.ch, https://orcid.org/0000-0003-4462-3173</li> <li>artin Knapmeyer, DLR Institute of Planetary Research, Rutherfordstr. 2, 12489 Berlin, ermany; martin.knapmeyer@dlr.de, https://orcid.org/0000-0003-0319-2514</li> <li>uénolé Orhand-Mainsant, Institut Supérieur de l'Aéronautique et de l'Espace SUPAERO, 10</li> </ul>
N ni A R G G C G C C	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>ikolaj Dahmen, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; kolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>nna Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748- 522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005 aris, France; kawamura@ipgp.fr</li> <li>mir Khan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, vitzerland; akhan@ethz.ch, https://orcid.org/0000-0003-4462-3173</li> <li>artin Knapmeyer, DLR Institute of Planetary Research, Rutherfordstr. 2, 12489 Berlin, ermany; martin.knapmeyer@dlr.de, https://orcid.org/0000-0003-0319-2514</li> <li>uénolé Orhand-Mainsant, Institut Supérieur de l'Aéronautique et de l'Espace SUPAERO, 10 /enue Edouard Belin, 31400 Toulouse, France; guenole.orhand-mainsant@isae-supaero.fr.</li> </ul>
N ni R G G M G G A ht	<ul> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>anstantinos.charalambous@imperial.ac.uk, https://orcid.org/0000-0002-9139-3895</li> <li>akolaj.dahmen@erdw.ethz.ch, https://orcid.org/0000-0002-9114-6747</li> <li>ana Horleston, School of Earth Sciences, University of Bristol, Wills Memorial Building, Queens bad, Bristol BS8 1RJ, UK; anna.horleston@bristol.ac.uk, https://orcid.org/0000-0002-6748-522</li> <li>aichi Kawamura, Université de Paris, Institut de physique du globe de Paris, CNRS, F-75005</li> <li>aris, France; kawamura@ipgp.fr</li> <li>mir Khan, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, vitzerland; akhan@ethz.ch, https://orcid.org/0000-0003-4462-3173</li> <li>artin Knapmeyer, DLR Institute of Planetary Research, Rutherfordstr. 2, 12489 Berlin, armany; martin.knapmeyer@dlr.de, https://orcid.org/0000-0003-0319-2514</li> <li>uénolé Orhand-Mainsant, Institut Supérieur de l'Aéronautique et de l'Espace SUPAERO, 10 /enue Edouard Belin, 31400 Toulouse, France; guenole.orhand-mainsant@isae-supaero.fr, tps://orcid.org/0000-0003-442-3173</li> </ul>

John-Robert Scholz, Max Planck Institute for Solar System Research, Justus-von-Liebig-Weg 3, 37077 Göttingen, Germany; <u>scholz@mps.mpg.de</u>, <u>https://orcid.org/0000-0003-1404-2335</u>
Fabian Euchner, Institute of Geophysics, ETH Zurich, Sonneggstr. 5, 8092 Zurich, Switzerland; fabian.euchner@sed.ethz.ch, <u>https://orcid.org/0000-0001-6340-7439</u>
W. Bruce Banerdt, Jet Propulsion Laboratory, California Institute of Technology, Pasadena, CA 91109, USA; <u>bruce.banerdt@ipl.nasa.gov, https://orcid.org/0000-0003-3125-1542</u>

### 802 Tables

Table 1. Number of seismic events for different event families and categories in the InSight v.5 catalog (until 2020/10/12; InSight Marsquake Service, 2021) that we use for magnitude calibration of four magnitude scales: spectral magnitude,  $M_{w,\text{spec}}^{\text{Ma}}$ , P- and S-body-wave magnitudes,  $m_b^{\text{Ma}}$  and  $m_{bS}^{\text{Ma}}$ , and 2.4 Hz magnitude,  $M_{2.4}^{\text{Ma}}$ . We only use events with distance estimates. Four events (S0183a, S205a, S0325a and S0357a, see Table S1) are excluded from the calibration of  $m_{bS}^{\text{Ma}}$ , because they fall into the low (S-wave) amplitude-zone of Mars (Giardini *et* al., 2020).  $M_{w,spec}^{Ma}$  is used as a calibration reference for the other scales for events with multiple amplitude measurements: it is most closely related to the moment magnitude  $M_W$ , less affected by noise, and can be determined for three out of five event categories.

	Event family	Event category	Number of events	MMa W,spec spectral domain	$m_{ m b}^{ m Ma}$ time domain	m <mark>Ma</mark> time domain	M <sup>Ma</sup> spectral or time domain
	v- ency	Low-frequency ( <i>LF</i> )	20	20	20	16	-
	Lov freque	Broadband ( <i>BB</i> )	9	9	9	8	-
	h-frequency	High-frequency ( <i>HF</i> )	50	50	-	-	50
		Very high- frequency (VF)	21	8	-	-	21
	Hig	2.4 Hz resonance (2.4Hz)	179	-	-	-	179

814	<b>Table 2.</b> Resulting magnitude scales for marsquakes. Spectral amplitudes $A_0$ (at 20 – 100 s) and
815	$A_{2.4,\text{spec}}$ (at 2 – 3 Hz) are in $m/\sqrt{Hz}$ , time domain amplitudes $A_P$ , $A_S$ (both at 0.2 - 0.5 Hz) and
816	$A_{24 \text{ pick}}$ (at 2 – 3 Hz) in <i>m</i> , and epicentral distance $\Delta$ in $\circ$ [degrees].

Name	Equation	Event family & categories	Calibrated Distance Range
Spectral magnitude	$M_{W,\text{spec},LF}^{\text{Ma}} = 2/3 \left( log_{10}(A_0) + (1.0 \pm 0.1) log_{10}(\Delta) + (12.6 \pm 0.5) \right)$	Low-frequency (LF, BB)	$25^{\circ} \le \Delta \le 100^{\circ}$
Spectral magnitude	$M_{W,\text{spec},HF}^{\text{Ma}} = 2/3(\log_{10}(A_0) + (0.8 \pm 0.1) \log_{10}(\Delta) + 12.8)$	High-frequency (HF, VF)	$3^{\circ} \leq \Delta \leq 30^{\circ}$
Body-wave magnitude (P-wave)	$m_{\rm b}^{\rm Ma} = \log_{10}(A_P) + 0.73 \log_{10}(\Delta) + 11.8$	Low-frequency ( <i>LF</i> , <i>BB</i> )	$25^{\circ} \le \Delta \le 100^{\circ}$
Body-wave magnitude (S- wave)	$m_{\rm bS}^{\rm Ma} = \log_{10}(A_{\rm S}) + 1.06 \log_{10}(\Delta) + 10.9$	Low-frequency ( <i>LF</i> , <i>BB</i> )	$25^{\circ} \le \Delta \le 35^{\circ}$ $60^{\circ} \le \Delta \le 100^{\circ}$
2.4Hz magnitude (time domain)	$M_{2.4,\text{pick}}^{Ma} = \log_{10}(A_{2.4,\text{pick}}) + 1.0 \log_{10}(\Delta) + 10.8$	High-frequency (VF, HF, 2.4Hz)	$3^{\circ} \leq \Delta \leq 35^{\circ}$
2.4Hz magnitude (spectral domain)	$M_{2.4,\text{spec}}^{Ma} = \log_{10}(A_{2.4,\text{spec}}) + 1.0 \log_{10}(\Delta) + 11.0$	High-frequency (VF, <i>HF</i> , 2.4Hz)	$3^{\circ} \leq \Delta \leq 35^{\circ}$

- 818 **Figure Captions**
- 819

Figure 1. Seismic wave propagation on Mars. The current interpretation of seismic signals on
Mars is that events in the high-frequency event family (*HF*, *VF*) excite guided waves in the crust or
in a shallow layer, while events in the low-frequency family (*LF*, *BB*) are located below the Moho;
these mantle waves are observed at long periods due to attenuation. A low-velocity layer (LVL)
could potentially explain the observed distant-dependent variation of S-wave amplitudes.
Modified after Giardini *et al.* (2020).

826

827 Figure 2. Amplitude computation from marsquake spectra and waveforms. (a) Theoretical 828 spectral shape of power spectral density (PSD) displacement spectra for events of the low-829 frequency (LF, BB; red) and high-frequency (HF, VF; blue) event families for various magnitudes. 830 Black line shows typical background noise. (b) Filtered displacement waveforms (vertical 831 component). Spectral amplitudes  $A_0$  and  $A_{2.4,spec}$  in (a) are determined from the long-period 832 plateau of the event displacement spectra and from a Lorentz curve fitted to the 2.4 Hz resonance 833 peak, respectively. Peak displacement amplitudes  $A_P$  and  $A_S$  in (b) are determined from the 834 filtered (2 - 6 s, top) displacement waveforms in the time window of the P- or S-wave arrivals. Note 835 that  $A_s$  is actually picked on the horizontals, while we only show vertical components for the sake 836 of simplicity. A<sub>2.4,pick</sub> in (b) is determined from the peak amplitude of the 2-3 Hz filtered 837 displacement waveform (bottom). Subplots in (c) and (d) show example spectra and waveforms 838 (vertical components) of representative marsquakes of each event category: *LF* – low-frequency, 839 BB - broadband, HF - high-frequency, VF - very high-frequency, and 2.4Hz - 2.4Hz ambient 840 resonance event.

841

**Figure 3. Magnitude calibration for marsquakes.** First, we calibrate the spectral magnitude  $M_{W,\text{spec}}^{\text{Ma}}$  from synthetic waveforms and theoretical considerations (see main text). Then we calibrate the other magnitude scales (body-wave magnitudes  $m_{b}^{\text{Ma}}$  and  $m_{bs}^{\text{Ma}}$ , and 2.4 Hz magnitude  $M_{2.4}^{\text{Ma}}$ ) relative to  $M_{W,\text{spec}}^{\text{Ma}}$  for events with multiple amplitude measurements ( $A_0$  - spectral plateau amplitude;  $A_P$ ,  $A_S$  - time domain peak P- and S-wave amplitudes at 0.2 – 0.5 Hz;  $A_{2.4}$ - resonance amplitude at 2 - 3 Hz in either time or spectral domain).

848

Figure 4. Marsquake magnitudes relative to spectral magnitude,  $M_{W,spec}^{Ma}$ . Body-wave magnitudes (a)  $m_b^{Ma}$  and (b)  $m_{bS}^{Ma}$  for low-frequency events (*LF*, *BB*), and 2.4Hz magnitudes (c)  $M_{2.4,pick}^{Ma}$  and (d)  $M_{2.4,spec}^{Ma}$  for high-frequency events (*HF*) relative to spectral magnitude,  $M_{W,spec}^{Ma}$ , which has been used as a calibration reference (and considered to scale with moment magnitude,  $M_W$ ). Markersize scales with event SNR. The calibration achieves good agreement across the different scales with standard deviations of  $\sigma = 0.3$  for  $m_b^{Ma}$ ,  $m_{bS}^{Ma}$ , and of  $\sigma = 0.2$  for  $M_{2.4,pick}^{Ma}$  and  $M_{2.4,spec}^{Ma}$ . The *HF* outlier in (c) is produced by S0490a.

857 Figure 5. Magnitude and distance distribution of marsquakes in the InSight catalog. Plot 858 shows all low-frequency (LF), broadband (BB), high-frequency (HF), very high-frequency (VF) and 2.4Hz resonance (2.4Hz) events. The preferred magnitude is  $M_W^{Ma} = M_{W,spec,LF}^{Ma}$  for LF and BB 859 events, to  $M_W^{Ma} = M_{W,spec,HF}^{Ma}$  for HF events, and to  $M_W^{Ma} = M_{2.4,spec}^{Ma}$  for VF and 2.4Hz resonance 860 861 events. Distances are derived from phase picks (InSight Marsquake Service, 2021), waveform 862 envelope alignment (Giardini *et al.*, 2020), and a simple crustal model (van Driel *et al.*, 2021), 863 respectively. Markersize scales with event SNR. Error bars show estimated magnitude and 864 distance uncertainties. Both distances and magnitudes for the most distant events are uncertain. 865

Figure 6. Comparison with magnitude scales on Earth. Comparison of (a)  $m_b^{Ma}$  and (b)  $m_{bs}^{Ma}$ 866 with the body-wave broadband magnitude, *mB* (Gutenberg, 1945a,b) for P- and S-waves on Earth 867

868 for all *LF* and *BB* events in the InSight v.5 catalog. Although developed for teleseismic waves at 2 -

869 20 s period on Earth, we are applying the mB relations here to Mars for  $A_P$  and  $A_S$  at 2 - 6 s period.

870

The good agreement of magnitudes on Mars and Earth suggests a similar attenuation of body-

871 waves at long periods and teleseismic distances.

### 873 Figures





Figure 1. Seismic wave propagation on Mars. The current interpretation of seismic signals on
Mars is that events in the high-frequency event family (*HF*, *VF*) excite guided waves in the crust or
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884 Figure 2. Amplitude computation from marsquake spectra and waveforms. (a) Theoretical 885 spectral shape of power spectral density (PSD) displacement spectra for events of the low-886 frequency (LF, BB; red) and high-frequency (HF, VF; blue) event families for various magnitudes. 887 Black line shows typical background noise. (b) Filtered displacement waveforms (vertical 888 component). Spectral amplitudes  $A_0$  and  $A_{2.4,spec}$  in (a) are determined from the long-period 889 plateau of the event displacement spectra and from a Lorentz curve fitted to the 2.4 Hz resonance 890 peak, respectively. Peak displacement amplitudes  $A_P$  and  $A_S$  in (b) are determined from the 891 filtered (2 - 6 s, top) displacement waveforms in the time window of the P- or S-wave arrivals. Note 892 that  $A_s$  is actually picked on the horizontals, while we only show vertical components for the sake 893 of simplicity.  $A_{2.4,\text{pick}}$  in (b) is determined from the peak amplitude of the 2-3 Hz filtered 894 displacement waveform (bottom). Subplots in (c) and (d) show example spectra and waveforms 895 (vertical components) of representative marsquakes of each event category: *LF* – low-frequency, 896 BB - broadband, HF - high-frequency, VF - very high-frequency, and 2.4Hz - 2.4Hz ambient 897 resonance event.



898 899

Figure 3. Magnitude calibration for marsquakes. First, we calibrate the spectral magnitude 900  $M_{W,\text{spec}}^{\text{Ma}}$  from synthetic waveforms and theoretical considerations (see main text). Then we calibrate the other magnitude scales (body-wave magnitudes  $m_{\rm b}^{\rm Ma}$  and  $m_{\rm bS}^{\rm Ma}$ , and 2.4 Hz magnitude 901  $M_{2.4}^{Ma}$ ) relative to  $M_{W,spec}^{Ma}$  for events with multiple amplitude measurements ( $A_0$  - spectral plateau 902 903 amplitude;  $A_P$ ,  $A_S$  - time domain peak P- and S-wave amplitudes at 0.2 – 0.5 Hz;  $A_{2.4}$ - resonance 904 amplitude at 2 - 3 Hz in either time or spectral domain).





907 **Figure 4. Marsquake magnitudes relative to spectral magnitude,**  $M_{W,spec}^{Ma}$ . Body-wave 908 magnitudes (a)  $m_b^{Ma}$  and (b)  $m_{bs}^{Ma}$  for low-frequency events (*LF*, *BB*), and 2.4Hz magnitudes (c) 909  $M_{2.4,pick}^{Ma}$  and (d)  $M_{2.4,spec}^{Ma}$  for high-frequency events (*HF*) relative to spectral magnitude,  $M_{W,spec}^{Ma}$ , 910 which has been used as a calibration reference (and considered to scale with moment 911 magnitude,  $M_W$ ). Markersize scales with event SNR. The calibration achieves good agreement 912 across the different scales with standard deviations of  $\sigma = 0.3$  for  $m_b^{Ma}$ ,  $m_{bs}^{Ma}$ , and of  $\sigma = 0.2$  for 913  $M_{2.4,pick}^{Ma}$  and  $M_{2.4,spec}^{Ma}$ . The *HF* outlier in (c) is produced by S0490a.

- 914
- 915



916  $\Delta$  [degrees] 917 **Figure 5. Magnitude and distance distribution of marsquakes in the InSight catalog.** Plot 918 shows all low-frequency (*LF*), broadband (*BB*), high-frequency (*HF*), very high-frequency (*VF*) and 919 2.4Hz resonance (2.4Hz) events. The preferred magnitude is  $M_W^{Ma} = M_{W,spec,LF}^{Ma}$  for *LF* and *BB* 920 events, to  $M_W^{Ma} = M_{W,spec,HF}^{Ma}$  for *HF* events, and to  $M_W^{Ma} = M_{2.4,spec}^{Ma}$  for *VF* and 2.4Hz resonance 921 events. Distances are derived from phase picks (InSight Marsquake Service, 2021), waveform 922 envelope alignment (Giardini *et al.*, 2020), and a simple crustal model (van Driel *et al.*, 2021), 923 respectively. Markersize scales with event SNR. Error bars show estimated magnitude and 924 distance uncertainties. Both distances and magnitudes for the most distant events are uncertain.



**Figure 6. Comparison with magnitude scales on Earth.** Comparison of (a)  $m_b^{Ma}$  and (b)  $m_{bs}^{Ma}$ with the body-wave broadband magnitude, mB (Gutenberg, 1945a,b) for P- and S-waves on Earth for all *LF* and *BB* events in the InSight v.5 catalog. Although developed for teleseismic waves at 2 -20 s period on Earth, we are applying the *mB* relations here to Mars for  $A_P$  and  $A_S$  at 2 - 6 s period. The good agreement of magnitudes on Mars and Earth suggests a similar attenuation of bodywaves at long periods and teleseismic distances.