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Crustal Anisotropy in the Martian Lowlands from Surface Waves

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- **Key Points:** 12
- Rayleigh and Love waves were detected on Mars; 13 • Seismic wave speed is directionally dependent over large-scales within Mars crust; 14 - The shear wave horizontal velocity is faster than the vertical velocity at $10 - 25 \,\mathrm{km}$ 15 depth in the lowlands between the event and the lander.

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17 Abstract

The largest seismic event ever recorded on Mars, with a moment magnitude of $4.7\pm$ 18 0.2, is the first event to produce both Love and Rayleigh wave signals. We measured their 19 group velocity dispersion between about 15 s and 40 s period and found that no isotropic 20 depth-dependent velocity model could explain the two types of waves wave simultane-21 ously, likely indicating the presence of seismic anisotropy. Inversions of Love and Rayleigh 22 waves yielded velocity models with horizontally polarized shear waves traveling faster 23 than vertically polarized shear waves in the top 10 km - 25 km. We discuss the possi-24 ble origins of this signal, including the preferred orientation of anisotropic crystals due 25 to shear deformation, alignment of cracks, layered intrusions due to an impact, horizon-26 tal layering due to the presence of a large-scale sediment layer on top of the crust, and 27 alternation of sedimentation and basalt layers deposits due to large volcanic eruptions. 28

29

Plain Language Summary

The largest marsquake recorded so far sent energy across the red planet in the form 30 of different kinds of waves, including two types of waves trapped near the surface. This 31 was the first time both types of so-called surface waves were detected on Mars. Com-32 bining measurements from these two types of surface waves allowed us to determine the 33 speed of other types of waves, i.e. waves that travel horizontally and that make rocks 34 move perpendicular to the direction of propagation. We found that these waves move 35 faster in the crust between 10 km and 25 km depth when the rocks oscillate in a direc-36 tion sub-parallel to the planet surface than if the rocks vibrate in the vertical direction. 37 This wave speed dependence can tell us about deformation mechanisms inside the crust. 38 We found that either an alternation of volcanic rocks and sediments layers due to vol-39 canic eruptions or internal layering within the crust due to an impact are the preferred 40 explanations for our observations. 41

42 **1** Introduction

The largest seismic event ever detected on Mars occurred on May 4 2022, and was recorded by the Seismic Experiment for Interior Structure (SEIS) Very Broadband (VBB) seismometer (Lognonné et al., 2020) of NASA's InSight mission (Banerdt et al., 2020). The event, labeled S1222a, occurred at 23:23:06.57 UTC (sol 1222) and has a moment

magnitude M_w estimated at 4.7±0.2 (Kawamura et al., 2022). This event is thus five 47 times larger than the second largest event recorded (Horleston et al., 2022) and had enough 48 energy to allow for the first direct observation of multi-orbiting Rayleigh waves on Mars 49 (Panning et al., 2022). Based on their calculated back azimuth of 101° (with a range of 50 $96^{\circ} - 112^{\circ}$) and epicentral distance $(37^{\circ} \pm 1.6^{\circ})$, the Mars Quake Service (MQS) deter-51 mined that this quality A (InSight Marsquake Service, 2020) event originated at 7.63° S, 52 170.67° E, near the North-South dichotomy, east of the landing site, and south of Cer-53 berus Fossae (Fig. 1). 54

Most remarkably, S1222a generated surface waves clearly observable on all three 55 components of the ground motion (Fig. 2). While Rayleigh waves have been previously 56 detected on Mars with impact events S1000a and S1094b (Kim et al., 2022), no Love waves 57 had been seen before. Since Rayleigh and Love waves are predominantly sensitive to dif-58 ferent elastic parameters (governing the speed of vertically and horizontally polarized 59 shear waves travelling horizontally, respectively), this gives us an opportunity to study 60 the presence of large-scale radial seismic anisotropy on Mars between the lander and the 61 event epicenter. Radial anisotropy, which is a type of transverse isotropy with a verti-62 cal symmetry axis, manifests as a difference in the velocities of vertically (V_{SV}) and hor-63 izontally (V_{SH}) polarized shear waves. It is quantified by parameter $\xi = V_{SH}^2/V_{SV}^2$ and 64 is a powerful tool to study the deformation history of a planet (e.g. Weiss et al. (1999)). 65

66 2 Data Analysis

We first removed the instrument response from the raw data to get the velocity record, 67 and rotated the U, V, and W components of the SEIS instrument to N, E, and Z. We 68 then rotated the N and E components to the radial (R) and tangential (T) coordinates 69 using a back azimuth of 101° as determined by the MQS (Kawamura et al., 2022). The 70 waveforms display clear wave packets on the Z and R components, characteristic of Rayleigh 71 waves, and a wave packet on the T component, identified as a Love wave (Fig. 2(a)). Us-72 ing two-pass, second-order Butterworth filters in different frequency bands (10 - 20 s,73 15 - 25 s, 20 - 30 s, 25 - 30 s, and 30 - 40 s), we found that S1222a generated surface 74 wave energy at periods between about 10 s and 40 s on all components of the seismo-75 gram (Fig. S1). 76

We subsequently carried out a multiple filter analysis (MFT) of the Z, R, and T 77 components to isolate the group velocities as a function of time (Dziewonski et al., 1969; 78 Herrmann, 2013). This consists in narrow bandpass filters using the operator $\exp\left[-\alpha(\omega-\omega_0)^2/\omega_0^2\right]$ 79 where ω_0 is the center frequency. Parameter α compromises between resolution in the 80 time and frequency domains. We measured the group velocities by picking the maximum 81 energy of the waveform envelope for each period (Fig. 2(b) - (d)). The velocity uncer-82 tainty is defined by a threshold energy of at least 1 dB of the energy maxima as in Beghein 83 et al. (2010). The MFT analysis shows that Rayleigh waves have peak energy between 84 about 15 s and 40 s period and group velocities increasing continuously up to about $\sim 3 \text{ km/s}$ 85 (Fig. 2(b) - (c)). The transverse component displays peak Love wave energy in the same 86 period range but with faster group velocities between about 3.2 km/s and 3.5 km/s (Fig. 2(d)). 87 We note that the T component MFT plot differs from the Z and R components in that 88 its dispersion curve first decreases between 10 s and 15 s and becomes relatively flat at 89 greater periods. 90

We tested that the effect of back azimuth uncertainties on the dispersion curves and found no significant change (Fig. S2). We also performed an ellipticity analysis in different period bands (Fig. S3). The elliptical and retrograde motion found between 15 s and 35 s confirm the Rayleigh wave nature of the signal.

We additionally observe discrepancies in the Rayleigh wave group velocities between 95 the Z and R components at periods >30 s, and even more so at periods >35 s for which 96 there is little overlap between the error bars (Fig. 2 and Fig. S3(a)). These differences 97 could be due to the lower waveform amplitude along the R component (Fig. S3(b)), mak-98 ing it more easily contaminated by noise. We considered stacking the R and Hilbert-transformed 99 Z components to improve the signal-to-noise ratio and obtain more reliable and consis-100 tent group velocities at periods >30 s. However, our ellipticity analysis showed that, while 101 at short period (15 - 30 s) the phase shift between the Z and R components is about 90° 102 (i.e., the major and minor axes of the ellipse are not parallel to the x- or y-axis), it is 103 not the case at 30-35 s, indicating that applying a Hilbert transform will not help. At 104 periods of 35-40 s, although the phase shift is close to 90° , the ellipticity (i.e., the ra-105 tio between Z and R) is much larger than that at short periods due to smaller ampli-106 tudes on the R component and may we not have a reliable signal. Because of these dis-107 crepancies between the Z and R components, we chose to invert measurements up to 30 s 108

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only. Inversions including all periods were performed for comparison. They are presented
 in the Supplementary Material and discussed in section 4.

The presence of transient signals, including glitches due to diurnal temperature variations on Mars, has the potential to contaminate the waveforms (e.g. Kim, Davis, et al. (2021)). We performed a glitch detection on the raw data using a synthetic glitch template (Scholz et al., 2020) and found no glitches near the first minor-arc surface wave arrivals (Fig. S4).

116 3 Method

The measured Love and Rayleigh wave dispersion curves were inverted jointly to 117 obtain one-dimensional (1-D) V_{SV} and V_{SH} models representing the average structure 118 between the station and the quake. The V_{SH} and V_{SV} model distributions were then com-119 bined to calculate radial anisotropy ξ . We used the Computer Programs in Seismology 120 (Herrmann, 2013) code package to generate the synthetic group velocity, and a Niche 121 Genetic Algorithm for the inversion (Koper et al., 1999; Li et al., 2021) together with 122 a B-spline parameterization, as described below. The results obtained with this approach, 123 hereafter referred to as Method 1, were compared against results from two other tech-124 niques (Method 2 and Method 3). Method 2 is a fully non-linear transdimensional hi-125 erarchical Bayesian approach using a reversible jump Monte Carlo Markov Chain (MCMC) 126 algorithm and a layered parameterization that is allowed to vary at each iteration (Weidner 127 et al., 2022). Method 3 also uses a MCMC approach but the parameterization is fixed 128 and uses B-splines. These methods are described in more details in the Supplementary 129 Material. 130

The period range at which the detected surface waves display strong energy (15 -30 s) corresponds to sensitivity to structure in the upper $\sim 70 - 80$ km, though it is important to note that Love wave sensitivity to V_{SH} is greatly reduced below 50 km (Fig. S5). Anisotropy was thus allowed in the models down to 80 km depth.

In the 0 – 80 km depth range, the V_{SV} and V_{SH} profiles were characterized by eight third-order B-splines. Since Rayleigh waves display some sensitivity down to about 130 km, we also inverted for isotropic shear wave velocities between 80 km and 130 km with a depth interval of 20 km for the depth anchor points. The prior constraint for all the shear wave velocities were uniform distributions from 1.5 km/s to 4.5 km/s. At depths greater than

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130 km, we used the InSight model KKS21GP, which was extracted from the models of 140 Khan et al. (2021) and Stähler et al. (2021) and selected to have a crust comparable to 141 that of Knapmeyer-Endrun et al. (2021). The V_P/V_S ratio was set to be 1.80 and the 142 density-to- V_S ratio was assumed to be 0.8. This choice was based on the ratios of the 143 receiver function study of Knapmeyer-Endrun et al. (2021), which are between 1.7 to 1.9 144 for V_P/V_S and 0.7 to 0.9 for density-to- V_S . In total, we searched for 200,000 models dur-145 ing the inversion, and the acceptable models were selected if their predicted group ve-146 locity at each period fits the measurements within uncertainties. The posterior distri-147 butions are based on these best data fitting models. 148

149 4 Results

We first inverted the Rayleigh wave dispersion curve and used the resulting 1-D 150 models to predict Love wave dispersion, and vice versa (Fig. S6). It is clear that isotropic 151 models obtained with the Rayleigh wave data are unable to explain the measured Love 152 wave dispersion and that the models obtained from Love waves alone cannot represent 153 the Rayleigh wave measurements. This same behavior has been widely observed in Earth's 154 upper mantle: Love waves usually have larger phase velocities than predicted by isotropic 155 models that fit Rayleigh wave phase velocities. It has also been detected in the conti-156 nental crust (e.g., Moschetti et al. (2010)). This apparent Love-Rayleigh discrepancy was 157 first observed on Earth by Anderson (1961) and can be resolved by introducing trans-158 verse isotropy in the model (Harkrider & Anderson, 1962). We thus performed joint Love 159 and Rayleigh wave inversions for V_{SH} and V_{SV} as described in section 3, and calculated 160 the resulting anisotropy ξ and Voigt V_S average model distributions: $V_S = \sqrt{(2V_{SV}^2 + V_{SH}^2)/3}$ 161 (Babuska & Cara, 2012). 162

Our average velocity models in the top 5 km and below 30 km depth exhibit sig-163 nificant uncertainties and variability between the different methods. The velocity obtained 164 with Method 1 at shallow depths is around 3.0 ± 0.1 km/s. V_S then increases to 3.2 ± 0.06 km/s 165 around 20 km depth and reaches 3.4 ± 0.1 km/s around 40 km below the surface. For 166 Method 2, the V_S models are characterized by about two layers with $V_S = 3.2 \pm 0.1$ km/s 167 in the top 20 – 30 km and V_S = 3.4 \pm 0.2 km/s below, down to 50 km depth. With 168 Method 3, V_S increases from 2.5 \pm 0.3 km/s to 3.0 \pm 0.75 km/s within the top 5 km, 169 $V_S = 3.4 \pm 0.8$ km/s at 20 km depth and $V_S = 3.5 \pm 0.2$ km/s at 40 km. 170

Our results also show a clear $V_{SH} > V_{SV}$ signal with $\xi \simeq 1.1$ between 5 km and 171 \sim 30 km depth (Fig. 3), consistent across all three methods within uncertainties (Figs. S7 172 and S8). The posterior model uncertainties do not, however, rule out $V_{SH} < V_{SV}$ in 173 the top 5 km and deeper than about 30 km. We additionally performed synthetic tests 174 that confirmed these results: In Fig. S9, we created and inverted an input model with 175 Method 1 that contained 10 % anisotropy in the 0 – 10 km depth range, 20 % at 10 – 30 km 176 depth, and 10% at 30-50 km. The model was isotropic at greater depths. The output 177 revealed that, while the positive anisotropy is well recovered in the 10-25 km depth range, 178 several solutions of opposite sign were found at shallower and greater depths. The lack 179 of resolution in the top $5 - 10 \,\mathrm{km}$ is likely due to the sensitivity of the data at the pe-180 riod range for which we have measurements. Shorter period data would be able to bet-181 ter constrain the shallower depths but have not been reliably recorded with the event 182 analyzed here. 183

Fig. 3 corresponds to inversions including only the Rayleigh wave data that were 184 consistent between the Z and R components, i.e. up to 30 s period. For completeness, 185 we also inverted the data measured on the Z and R separately up to 40 s (Fig. S10). These 186 tests show that including periods above 30 s yields model uncertainties that depend on 187 which component the Rayleigh waves are measured. In addition, while most models still 188 display $\xi > 1$ between 10 km and 25 km, the error bars are significantly larger. This 189 dependence of the models on the data at periods greater than 30 s validates our deci-190 sion to only present results based on the shorter period data in Fig. 3. 191

¹⁹² 5 Discussion

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5.1 Seismic Velocities

Lognonné et al. (2020) discussed the rock types, including sedimentary, impact-related, and volcanic, and their associated seismic velocities in the Martian crust around the In-Sight landing site. They found that the low velocities they had obtained for the shallowest crustal layer (upper ~8 km) suggest altered and/or damaged upper volcanic rocks. The velocities in the underlying layer (around 8 – 20 km depth), with $V_P = 4.5 - 6$ km/s and $V_S = 2.5 - 3.3$ km/s, were consistent with altered, iron-rich basaltic rocks. Similar velocities were later found by Knapmeyer-Endrun et al. (2021).

In our study, Methods 1 and 2 predict $V_S = 3.0 \pm 0.1$ km/s and $V_S = 3.2 \pm 0.1$ km/s, 201 respectively, in the top 5 - 10 km, which is larger than the 1.7 - 2.1 km/s estimated by 202 Lognonné et al. (2020) and Knapmeyer-Endrun et al. (2021) at the landing site. Method 203 3 gives lower V_S values between 2.5 \pm 0.3 km/s and 3.0 \pm 0.75 km/s in that depth 204 range, but they are still larger than found beneath the lander. Using $V_P/V_S = 1.8$ (Knapmeyer-205 Endrun et al., 2021), we calculated that this corresponds to $V_P = 4.5 - 5.4$ km/s, which 206 is still compatible with damaged or altered basalt, but also with impact melt, impact 207 breccia, and sandstone (Lognonné et al., 2020). 208

Between ~10 km and 30 km depth, V_S is consistent among the three methods with values around 3.1 - 3.5 km/s. Using a $V_P/V_S = 1.8$ (Knapmeyer-Endrun et al., 2021), this correspond to $V_P \simeq 5.6 - 6.3$ km/s. Compared to Lognonné et al. (2020), our velocity models at those depths are thus compatible with volcanic rocks (damaged, altered, or compact basalt) and with impact melt.

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5.2 Seismic Anisotropy

The presence of seismic anisotropy, i.e. the directional dependence of seismic wave velocity, is well documented on Earth (Dziewonski & Anderson, 1981; Montagner, 1994). In Earth's upper mantle, it is considered a strong indicator of past (in the mantle lithosphere) or present (in the asthenosphere) strain coherent over large scales, resulting in textural changes such as lattice or crystal preferred orientation (LPO or CPO) of intrinsically anisotropic olivine crystals due to mantle deformation (Nishimura & Forsyth, 1989; Karato & Wu, 1993; Montagner, 1998; Becker et al., 2003; Yuan & Romanowicz, 2010).

In Earth's crust, seismic anisotropy is often attributed to the shape preferred ori-222 entation (SPO) of isotropic material with contrasting elastic properties (Backus, 1962), 223 such as the preferred alignment of melt pockets or lenses (Bastow et al., 2010), the dis-224 tribution of fluid-filled cracks, or preferentially oriented pore spaces (Crampin, 1981; Crampin 225 & Booth, 1985; Crampin, 1989). SPO can be caused by sedimentary layering, magmatic 226 intrusions, or the presence of faults. In this case, the observed large-scale seismic anisotropy 227 represents the effective medium comprised of unmapped fine (subwavelength thickness) 228 isotropic layered structures (Backus, 1962). Alternatively, observations of anisotropy with 229 $V_{SH} > V_{SV}$ in Earth's mid-to-lower continental crust are often attributed to the align-230 ment of minerals such as mica in ductile flow (e.g., Brownlee et al. (2017)). In Tibet, for 231

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instance, it has been proposed that this type of anisotropy could be due to underthrusting of the Indian crust beneath Tibet or could result from ongoing crustal thinning (Shapiro
et al., 2004).

Our results indicate a robust seismic anisotropy signal in the 10 – 30 km depth range (Fig. 3). Constraints from gravity data show that the crustal thickness along the eventstation path, which samples the lowlands, is around 40 km (Wieczorek et al., 2022). In addition, recent P- and S-wave receiver function analyses with seismic data collected by SEIS have shown that the crustal thickness at the lander site is between 20 km and 40 km (Knapmeyer-Endrun et al., 2021), with a preference for 40 km with PP receiver functions (Kim, Lekić, et al., 2021).

In order to observe seismic anisotropy, there needs to be coherent preferred orientation over large distances. Specifically, producing $V_{SH} > V_{SV}$ would require either the preferential horizontal alignment of fast axes of anisotropic minerals or cracks. Therefore, we rule out the alignment of cracks (Crampin, 1981; Crampin & Booth, 1985; Crampin, 1989) as a likely explanation for the observed signal since alignment of vertically-oriented cracks would yield $V_{SV} > V_{SH}$. In addition, at depths greater than about 20 km, the pressure is such that cracks tend to close.

Another possibility is horizontal layering of isotropic material (with layer thicknesses 249 much smaller than the dominant wavelength), i.e. extrinsic or apparent anisotropy. Where 250 the different layers have contrasting elastic properties, an overall anisotropic medium re-251 sults. In order to assess the degree of heterogeneity required, we have performed some 252 simple calculations using Backus averaging (Backus, 1962). Fig. S11 shows the result-253 ing ξ parameter for a range of degrees of shear wave velocity perturbations, and layer 254 thickness ratios. Producing an anisotropic signal compatible with our observations re-255 quires a relative velocity perturbation of $\sim 20\%$, when fast and slow layer thicknesses 256 are approximately equal. As either material begins to dominate the ensemble, the re-257 quired perturbation significantly increases. Such a mechanism is observed in sedimen-258 tary basins on Earth (Sayers, 1998). On Mars, there is evidence for past floods and sed-259 imentary transportation near the rim at the Pathfinder landing site (Rodriguez et al., 260 2019), as well as the presence of rivers and even an ocean that may have had enough wa-261 ter to fill the Vastitas Borealis, the Northern basin in the upper hemisphere (Carr & Head III, 262 2003; Di Achille & Hynek, 2010; Villanueva et al., 2015). Such a standing body of wa-263

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ter could have deposited sediments over large distances (Head et al., 2001). However, 264 McLennan (2012) estimated the global average thickness of the sediment layer to be less 265 than 2 km. Although the sedimentary rocks at several locations can reach greater depths, 266 they are shallower than 10 km in many places (e.g., 5 km in Gale crater (Grotzinger et 267 al., 2015), 3 – 6 km in Juventae Chasma (Grotzinger & Milliken, 2012), and 3 km in Med-268 uase Fossae (Bradley et al., 2002). Therefore, the observed anisotropy at depth between 269 10 km and 30 km is unlikely due to sedimentary rocks alone unless they extend to much 270 greater depths than expected. 271

Layered sill complexes have been suggested to produce strong (> 8 %) positive ra-272 dial anisotropy at 5 - 18 km depth below silic calderas (Jiang et al., 2018), but it is usu-273 ally very localized in depth and spatial extent. The strength and extent of the anisotropy 274 observed here would require a lot of melt distributed in the Martian crust over a broad 275 depth range. Alternatively, with Mars hosting the largest volcano of the solar system north-276 east of InSight, Olympus Mons, as well as Elysium Mons to the north, multiple volcanic 277 events over broad scales could have led to an alternation of basalt layers deposit and sed-278 imentation, which could potentially have generated radial anisotropy, though, again, it 279 would have to extend over rather large depths. Basalt packages in a terrestrial context 280 (for example, those associated with large igneous provinces, e.g., Planke (1994); Christie 281 et al. (2006)) show a high degree of velocity layering, variation within which easily ex-282 ceeds the 20 % perturbation required to reproduce our inferred ξ value (Baan et al., 2007). 283 Much of this heterogeneity is associated with variations of porosity and sub-aerial frac-284 turing (Planke, 1994) which might be expected to be reduced at greater depth in the Earth. 285 It has been suggested, however, that on Mars pore space might persist to the observed 286 seismic discontinuity at 20 km below the InSight lander (Knapmeyer-Endrun et al., 2021; 287 Wieczorek et al., 2022). 288

Another option is the presence of layered intrusions due to a single or multiple im-289 pacts (Wilhelms & Squyres, 1984; Nimmo & Tanaka, 2005; Andrews-Hanna et al., 2008; 290 Marinova et al., 2008). It has been proposed that the crustal dichotomy, characterized 291 by a difference of up to 30 km in crustal thickness between the northern lowlands and 292 the southern highlands, corresponds to a crater formed by a large impact. Such an im-293 pact could have contributed to the evolution of the proto-crust on Mars and other ter-294 restrial planets, and could have caused internal layering within the crust. This has been 295 shown to be the case on Earth: Latypov et al. (2019) studied the exposed impact melt 296

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sheet of the Sudbury Igneous Complex located in Canada and showed that such ancient
 and compositionally primitive Hadean impact melt sheets underwent large-scale igneous
 differentiation.

This scenario is also consistent with the results of Li, Beghein, Lognonné, et al. (2022) 300 who constrained crustal V_{SV} in the highlands by measuring and inverting the dispersion 301 of minor-arc (R1) and major-arc (R2) Rayleigh waves, i.e. the wave that travels the short-302 est and the longest distance, respectively, along the great circle path between the epi-303 center and the instrument, according to ray theory. They found V_{SV} between about 3.3 km/s 304 and 3.5 km/s in the highlands, which is greater than the crustal V_{SV} of 3.0 - 3.2 km/s 305 found here in the lowlands, but similar to the 3.4 km/s we found for crustal V_{SH} in the 306 low lands (Fig. S8(b), S8(e), and S8(h)). This similarity in velocity between V_{SH} in the 307 lowlands and V_{SV} in the highlands could indicate that the average Martian upper crustal 308 shear wave velocities might have been originally similar in the lowlands and highlands. 309 The impact-related horizontal melt sheets created at depth in the lowlands could have 310 lowered the shear wave velocity in the vertical direction (i.e., V_{SV}) but left V_{SH} unchanged 311 (Fig. S12(a)). 312

Finally, we cannot entirely rule out the lattice preferred orientation of anisotropic 313 crystals in the Martian crust to explain the observed signal, as proposed for Earth's mid-314 to-lower continental crust (Shapiro et al., 2004). There is currently no large underthrust-315 ing on Mars such as that seen in Tibet (Shapiro et al., 2004), and there is no evidence 316 of large-scale crustal thinning between the lander and the event epicenter, but one could 317 envision that anisotropic crystal fabrics could be preserved in the crust from past defor-318 mation such as an impact. If the impact melted the crust and it began to flow, fabric 319 could have been generated over relatively large-scales and become frozen-in as rocks cooled 320 down. 321

We end this discussion by pointing out that, importantly, the observed anisotropy ($\xi > 1$) in this study does not contradict the previously observed anisotropy beneath the InSight lander site (Li, Beghein, Wookey, et al., 2022). First, the anisotropy in Li, Beghein, Wookey, et al. (2022) reflects the local structure (i.e., within 30 km) beneath the lander, and this study represents the average structure along the > 2,000 km long path between the epicenter and the station. Second, the anisotropy in Li, Beghein, Wookey, et al. (2022) was restricted to the topmost layer (around 8 km), whereas the most re-

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liably resolvable depths in the present study are 10-25 km. Last, but importantly, the 329 direction of the ray path is near-vertical for the body waves in Li, Beghein, Wookey, et 330 al. (2022) and is horizontal for the surface waves in this study. Therefore, the shear-wave 331 velocity we constrain from the surface waves is that of a horizontally propagating shear-332 wave. The definitions of V_{SV} are not thus the same in both studies since the directions 333 of the polarization of the SV waves (perpendicular to the direction of the seismic rays) 334 are different (Fig. S12(b) – (c)). Subsequently, the $V_{SV} < V_{SH}$ anisotropy in this study 335 does not contradict the $V_{SV} > V_{SH}$ anisotropy in Li, Beghein, Wookey, et al. (2022) 336 and they reflect different properties of the materials in the Martian crust. In fact, the 337 modeling done by Li, Beghein, Wookey, et al. (2022) to explain their observations used 338 a horizontally transverse isotropic medium, which reflects the presence of azimuthal anisotropy. 339 Here, the signal reflects radial anisotropy, also called polarization anisotropy or verti-340 cally transverse isotropy (VTI), which has hexagonal symmetry with a vertical symme-341 try axis. The observations beneath the lander require a medium with a higher degree 342 of symmetry. 343

6 Conclusion

We analyzed the Rayleigh and Love waves generated by the largest event ever recorded 345 on Mars. The waveforms contain surface wave energy on all three components of the seis-346 mogram at periods between 15 s and 40 s. We measured the group velocity dispersion 347 of these waves on the vertical, radial, and transverse components of the seismogram us-348 ing a multiple filter analysis. We found that Rayleigh waves group velocities increase con-349 tinuously up to about 3 km/s between 15 s and 40 s period. Love wave group velocities 350 are systematically larger, between about 3.2 km/s and 3.5 km/s in the 15 - 40 s period 351 range. Because of a discrepancy between the Rayleigh wave group velocities in the trans-352 verse and vertical directions at periods greater than 30 s, we only used data up to 30 s. 353

Isotropic inversions of the Love wave data were unable to explain the Rayleigh wave data and vice versa. The Love and Rayleigh wave dispersion curves were thus inverted jointly for shear wave velocities and radial anisotropy. We found a well-resolved signal in the crust, with $V_{SH} > V_{SV}$ ($\xi > 1$) between 10 km and 25 km depth.

LPO of anisotropic crystals inside the crust is difficult to reconcile with our current knowledge of the martian crust as it would require a large-scale shear mechanism

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- ³⁶⁰ such as crustal thinning or underthrusting. However, we could be seeing frozen-in fab-
- ric due to melting followed by flow (and thus LPO) in the crust after an impact, for in-
- 362 stance. The observed anisotropy may alternatively be apparent rather than intrinsic. We
- rule out cracks as a possible explanation for this signal since vertical cracks would yield
- $V_{SH} < V_{SV}$. Horizontal layering of sediments is unlikely to explain the signal as well
- ³⁶⁵ because they would not extend to great enough depths. Our preferred explanations in-
- clude an alternation of isotropic basalt layers deposit and sedimentation caused by vol-
- canic eruptions, or the presence of isotropic layered intrusions due to an impact.



Figure 1. Location of event S1222a (red star) and SEIS (black triangle). The grey shaded area represents the uncertainties in back azimuth and epicentral distance and the three black curves are possible great circle paths of the surface waves.



Figure 2. (a) Deglitched waveforms on the to Z, R, T coordinates. A second order, zero-phase shift bandpass filter was applied between 10 s and 60 s period. (b)-(d) Multiple filter analysis of the vertical (a), radial (b), and tangential (c) components of the VBB.



Figure 3. V_S (a) and ξ (a) models resulting from Method 1 applied to the measured dispersion curves. The vertical black dashed line in (b) represents an isotropic model for reference. The red horizontal bars are the uncertainties calculated by averaging the variance obtained by each of the three methods.

³⁶⁸ 7 Open Research

Data and materials availability: The InSight event catalogue (Clinton et al., 2021; Ceylan et al., 2022; InSight Marsquake Service, 2022a, 2022b) and waveform data are available from the IRIS-DMC, NASA-PDS (InSight Mars SEIS Data Service, 2019a) and IPGP data center (InSight Mars SEIS Data Service, 2019b).

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